

Precision primordial nucleosynthesis

(with the new code “PRIMAT”)

Alain Coc

(Centre de Sciences Nucléaires et de Sciences de la Matière, Orsay)

Cyril Pitrou, Elisabeth Vangioni, Jean-Philippe Uzan

(Institut d'Astrophysique de Paris)

The three observational evidences for the Big Bang Model

1. The expansion of the Universe

Galaxies move away from each other according to Hubble's law:

$V = H_0 \times D$ with $H_0 \approx 70 \text{ km/s/Mpc}$, the Hubble parameter (or “constant”).

More precisely distances $\propto a(t)$, the cosmological scale factor

2. The Cosmic Microwave Background radiation (CMB)

A black body radiation at 2.7 K corresponding to the redshifted spectrum emitted when the universe became transparent

3. Primordial nucleosynthesis

Reproduces the “light-elements” (^4He , ^2H or D , ^3He and ^7Li) primordial abundances over a range of nine orders of magnitudes.

Big Bang Nucleosynthesis probe of new physics

- ❑ First determination of the baryonic density of the Universe, $(1-3) \times 10^{-31} \text{g/cm}^3$ [*Wagoner 1973*], need for baryonic dark matter
 - Baryonic density $\rho_B \approx 4.5 \times 10^{-31} \text{g/cm}^3$ from the anisotropies in the Cosmic Microwave Background radiation,
- ❑ First determination of the number of light neutrino families, $N_\nu \leq 3$ [*Yang, Schramm, Steigman, Rood 1979*]
 - Number of neutrino families $N_\nu = 2.984 \pm 0.008$ [*LEP experiments*]
- ❑ Today, still extensively used to constraint physics beyond the Standard Model. See e.g. review by *Iocco et al. 2009*

Determination of primordial abundances

Primordial abundances :

1) Observe a set of primitive objects born when the Universe was young

- ^4He in H II (ionized H) regions of blue compact galaxies
- ^3He in H II regions of *our* Galaxy
- D in remote **cosmological clouds** (i.e. at high redshift) on the line of sight of quasars
- ^7Li at the surface of low metallicity* stars in the halo of our Galaxy

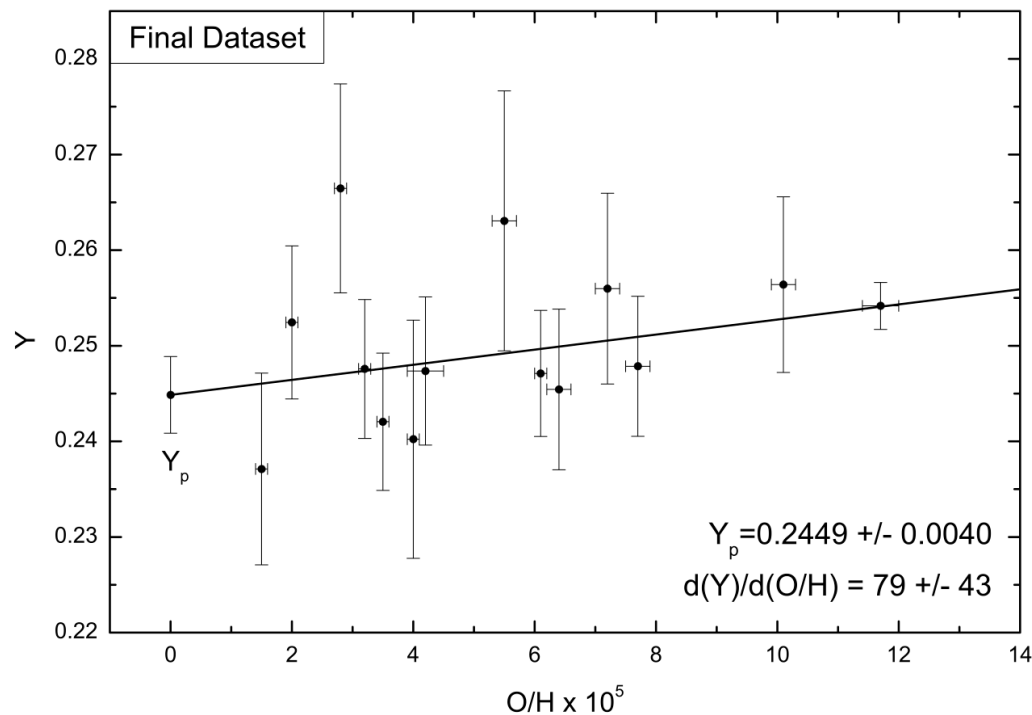
2) Extrapolate to zero metallicity* : $\text{Fe}/\text{H}, \text{O}/\text{H}, \text{Si}/\text{H}, \dots \rightarrow 0$

*In astrophysics: “metals” = everything beyond helium

Notation : $[\text{X}/\text{H}] \equiv \log(\text{X}/\text{H}) - \log(\text{X}_{\odot}/\text{H}_{\odot})$, $\text{X}=\text{Fe}, \text{O}, \dots$

^4He observations in blue compact galaxies

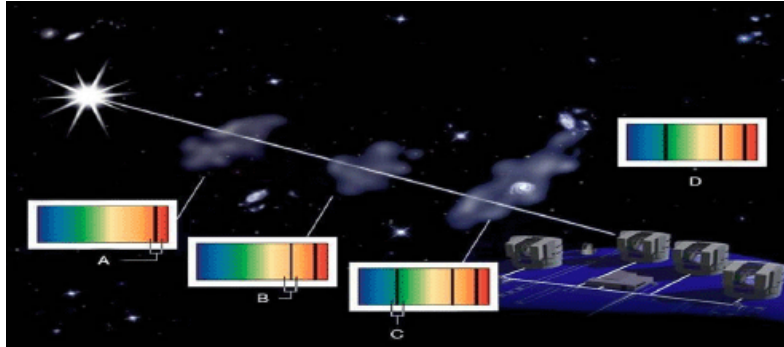
Observations: ^4He from a sample of 86 H II regions in 77 blue compact galaxies [*Izotov, Thuan & Stasinska 2007*] with additional infrared line [*Izotov, Thuan & Guseva 2014*]



$Y_p = ^4\text{He}$ mass fraction

Analysis : with new atomic and collisional emission data, He I emissivity and IR line included $Y_p = 0.2449 \pm 0.0040$ (^4He mass fraction) [*Aver, Olive, Porter & Skillman 2015*] (1.6% precision)

D/H observations in a cosmological cloud

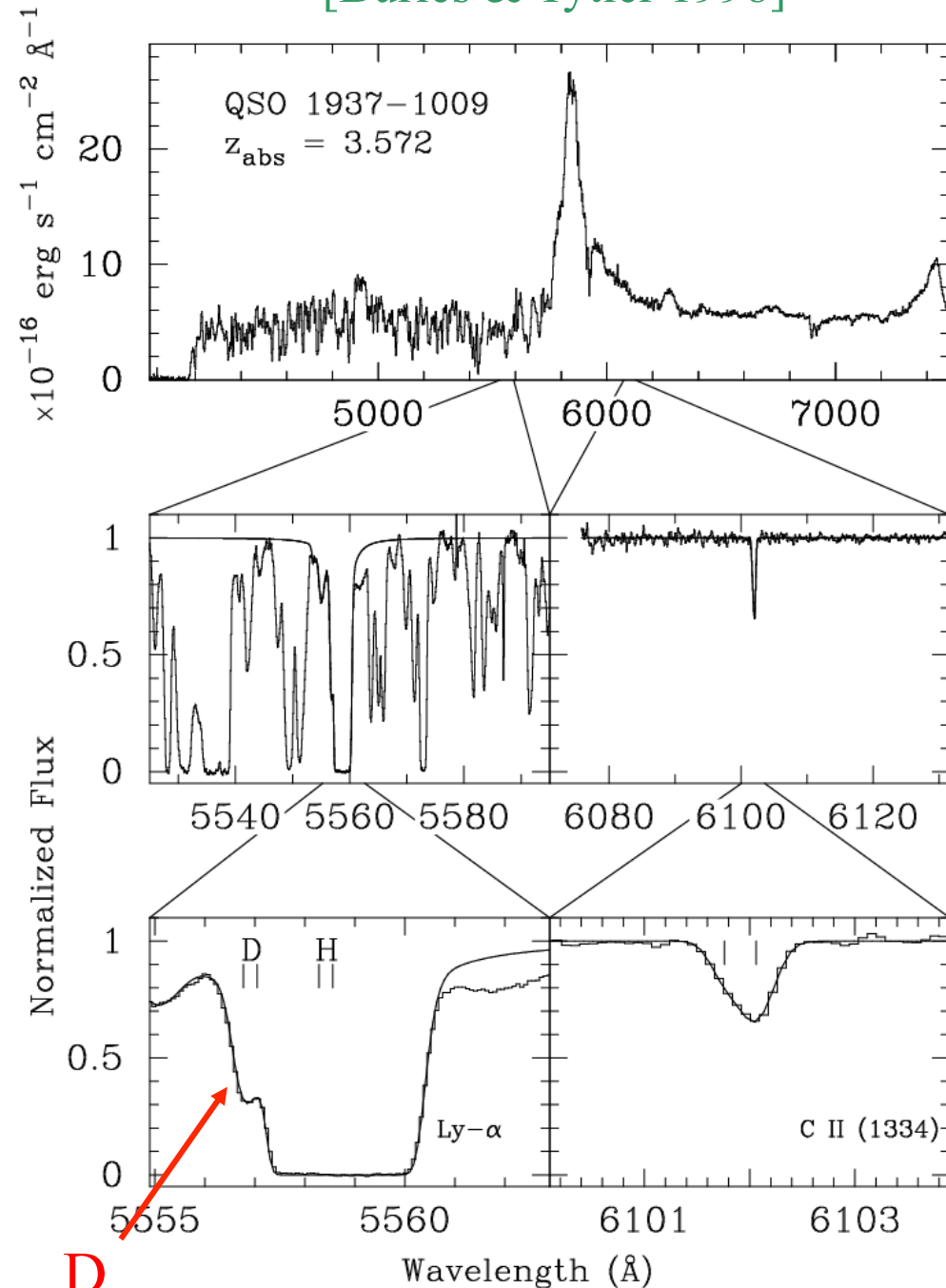


Cloud at redshift of $z = 3.6$
on the line of sight of
quasar QSO 1937-1009

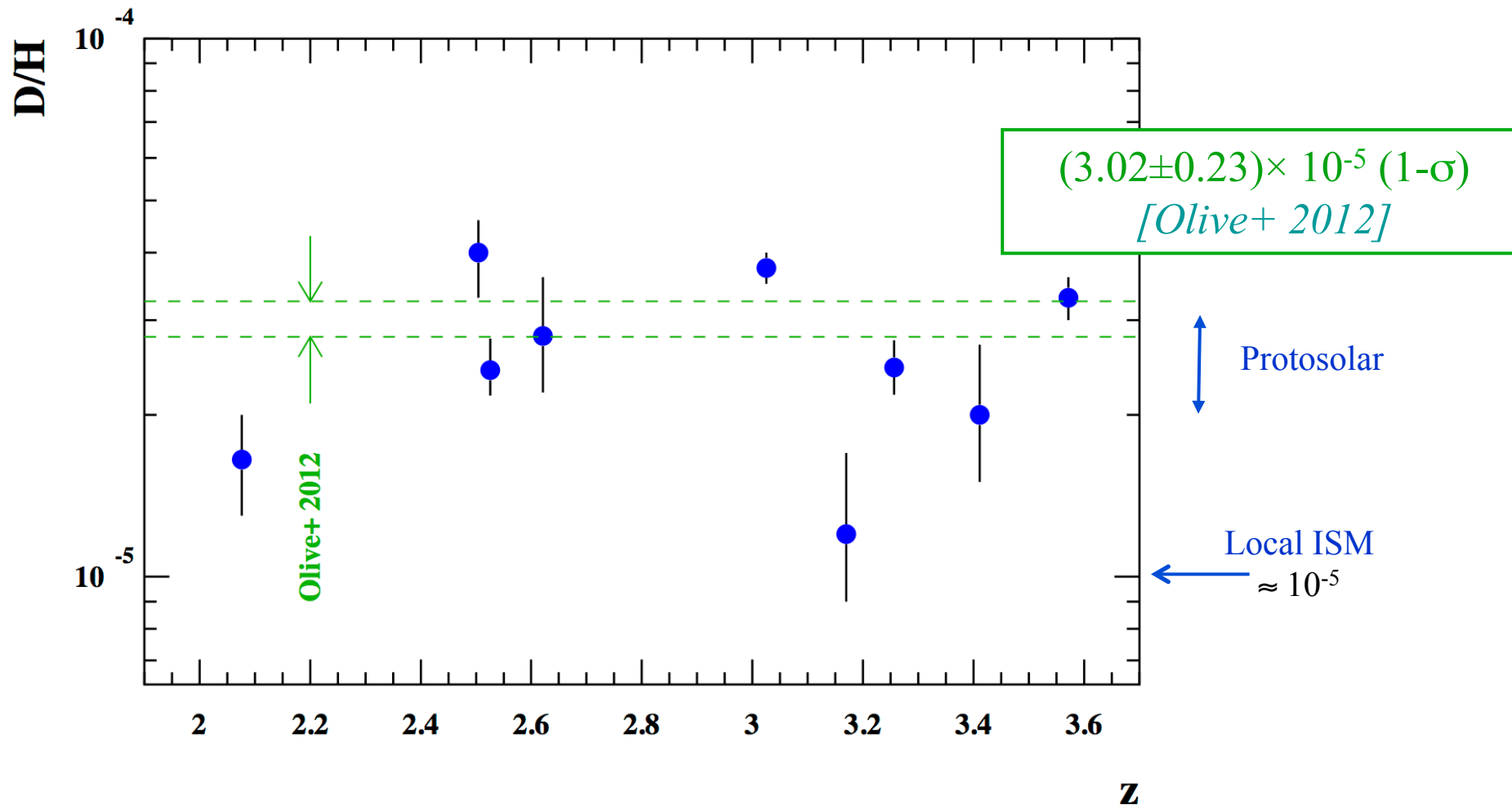
Observations :

- D/H ratio at high redshift
from the depth/width of
absorption lines

[Burles & Tytler 1998]

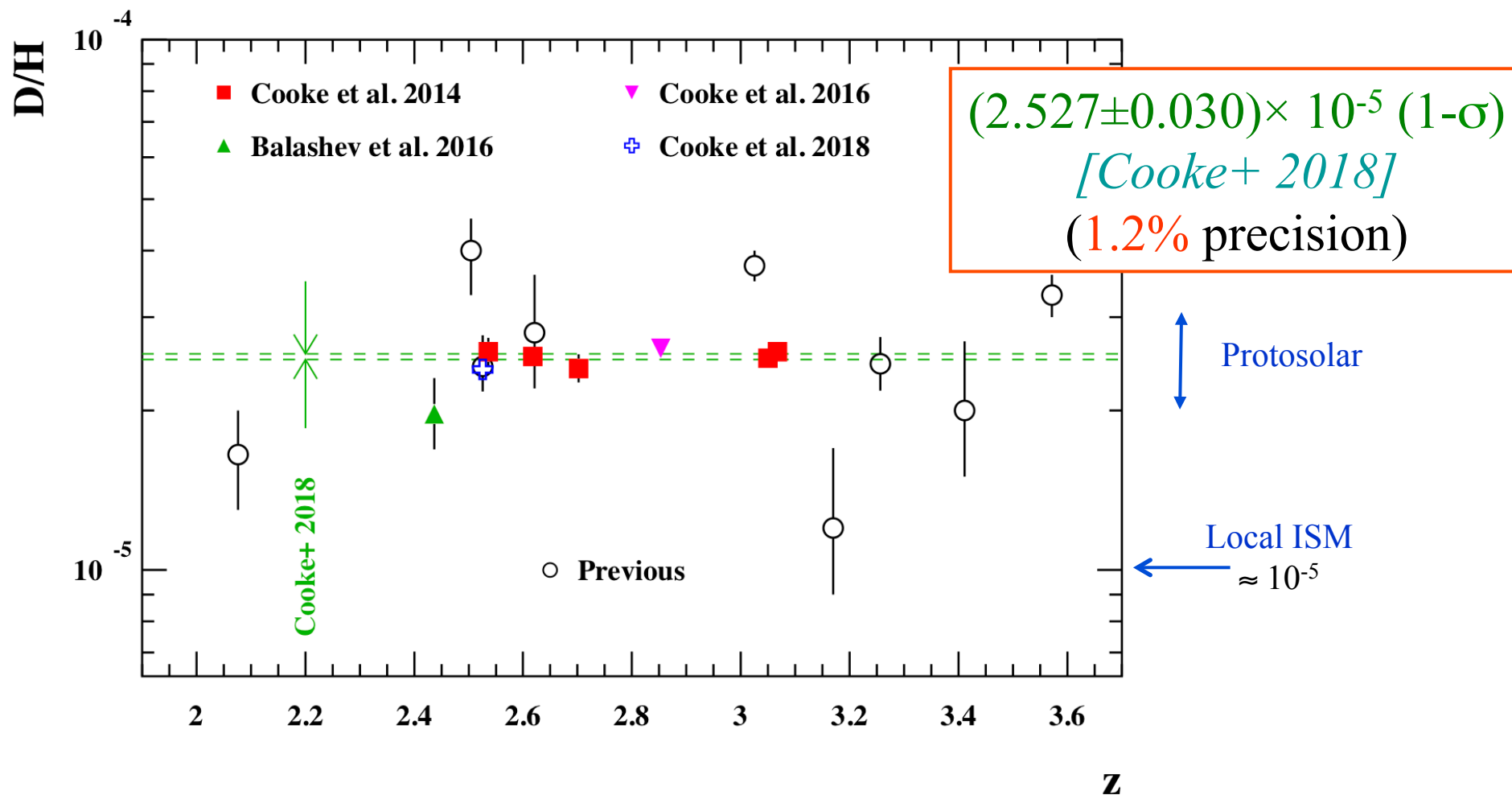


D/H observations in cosmological clouds



Burles & Tytler 1998; O'Meara+ 2001, 2006; Pettini+ 2001, 2008, 2012; Kirkman+ 2003, Crighton+ 2004; Srianand+ 2010; Cooke+ 2011; Fumagalli+ 2011

D/H observations in cosmological clouds

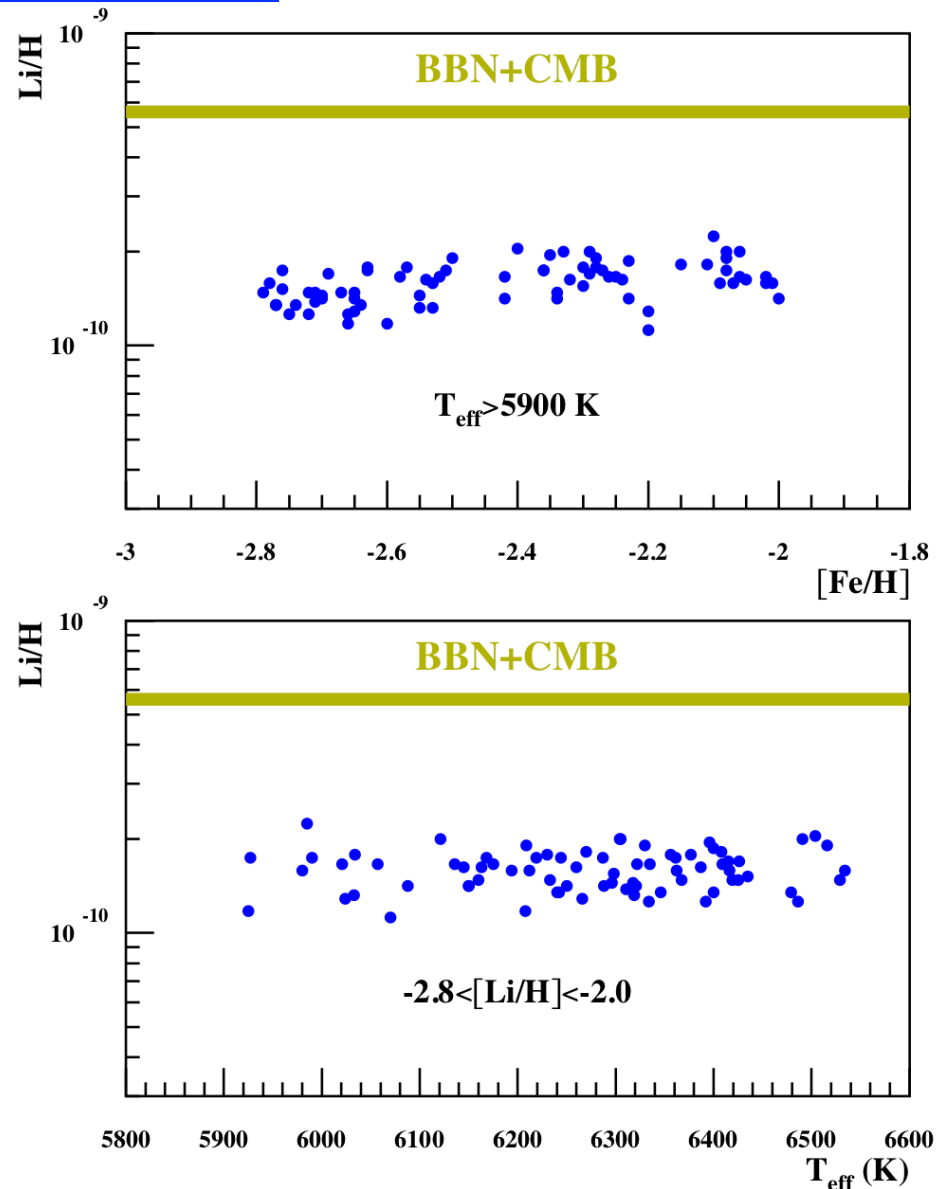


Burles & Tytler 1998; O'Meara+ 2001, 2006; Pettini+ 2001, 2008, 2012; Kirkman+ 2003, Crighton+ 2004; Srianand+ 2010; Cooke+ 2011; Fumagalli+ 2011; Cooke+ 2016; 2018

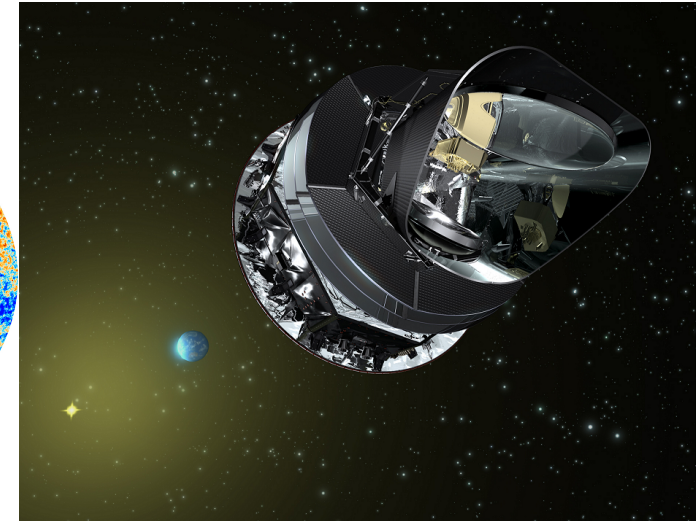
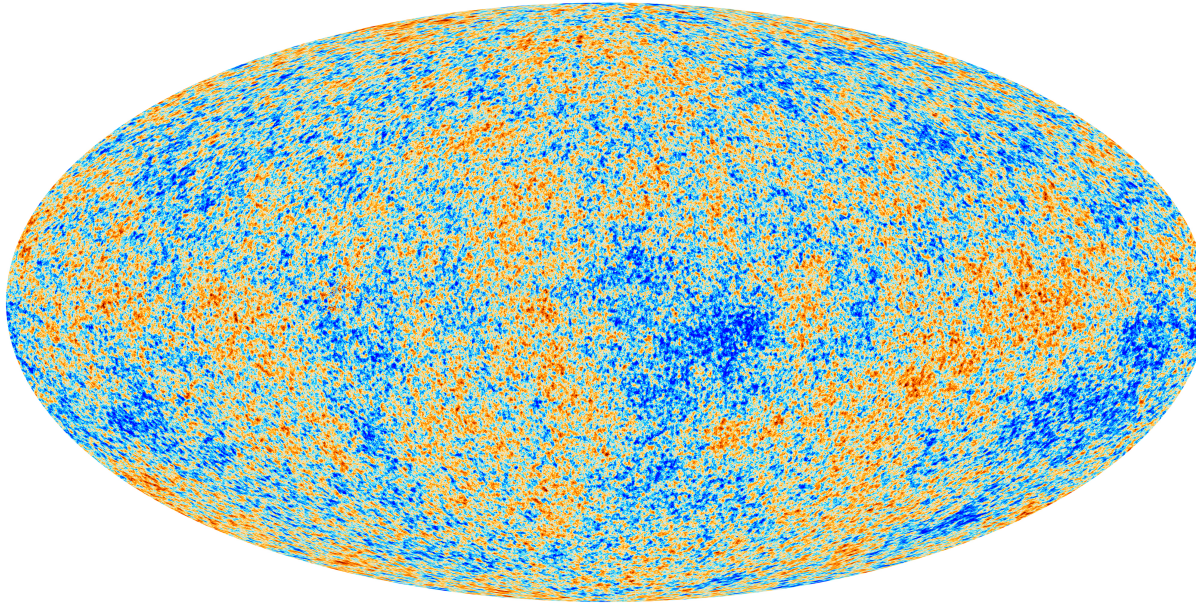
Primordial Li from observations

[M. Spite, priv. comm.]

- Oldest (low metallicity) stars in galactic halo
- Lifetime of $M < 0.9 M_{\odot}$ stars > 15 Gy
- For $T_{\text{eff}} > 5900$ K, no deep convection and no Li surface depletion (?)
- $\text{Li}/\text{H} = (1.58 \pm 0.35) \times 10^{-10}$
[Sbordone et al. 2010]



Density components of the Universe



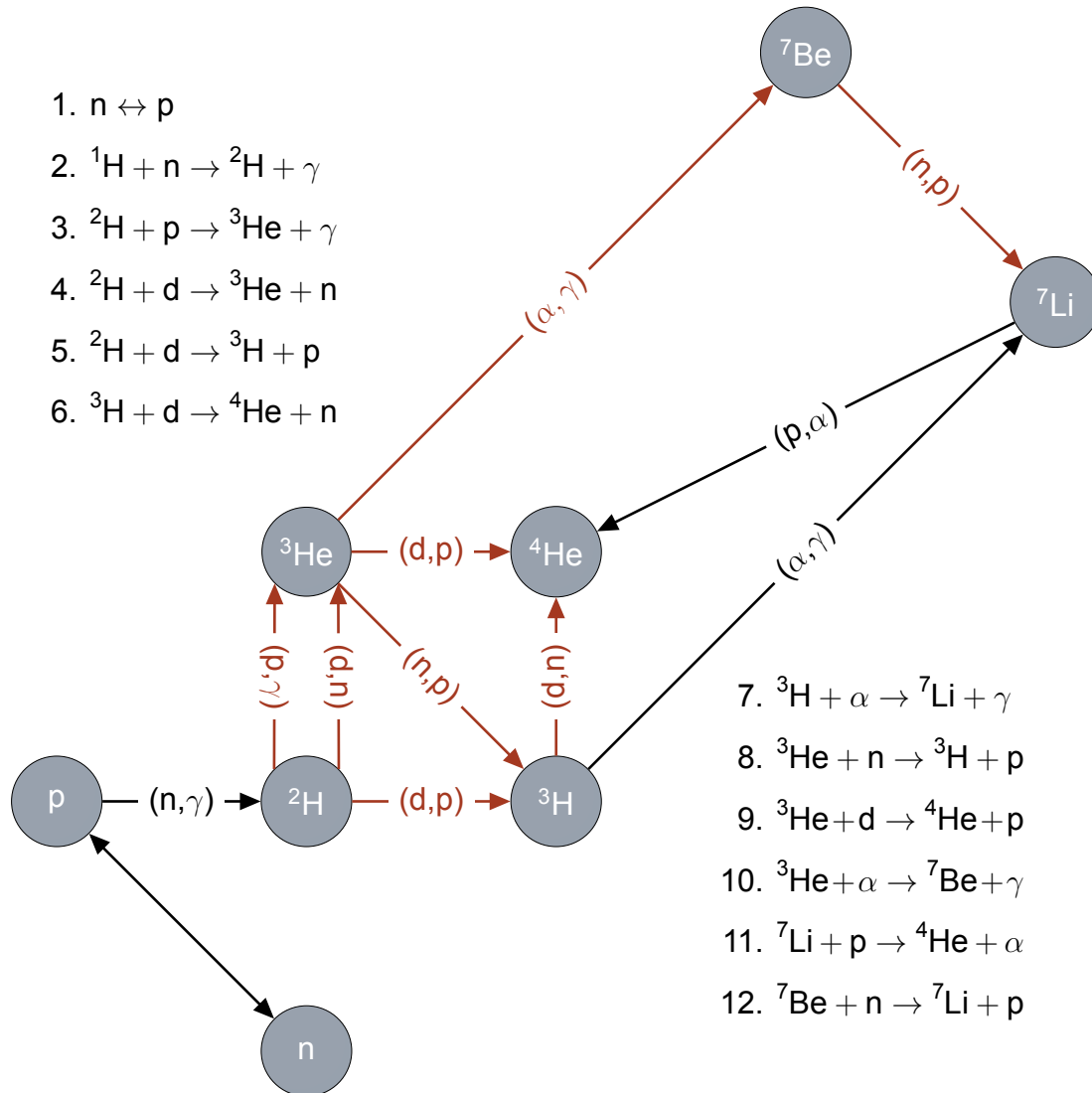
$$\Omega \equiv \frac{\rho}{\rho_{\text{Critical}}}$$

$$h = H_0 / 100 \text{ km/s/Mpc} \\ \approx 0.6727 \pm 0.0066$$

Some Ω values [<i>Ade+ 2016 (Planck)</i>]		
Radiation (CMB)	Ω_R	$5 \cdot 10^{-5}$
Visible matter	Ω_L	≈ 0.003
Baryons	Ω_b	0.049
Dark Matter	Ω_c	0.264
Vacuum	Ω_Λ	0.688
Total	Ω_T	≈ 1.0

$$\Omega_b h^2 = 0.02225 \pm 0.00016 \\ (\text{0.7\% precision}) \\ [\textit{Ade+ 2016 (Planck)}]$$

Big Bang Nucleosynthesis calculations



Needs:

- Reaction rates (~ 400 to reach CNO)
- Density $\rho_b(t)$, ions and photons $T(t)$ and neutrino $T_\nu(t)$ temperatures as a function of time

New public code

PRIMAT [Pitrou+ 2018]

Dynamics of the expanding Universe

Einstein equation & Friedmann-Lemaître-Robertson-Walker metrics

$$ds^2 = g_{\alpha\beta} dx^\alpha dx^\beta = dt^2 - a^2(t) \left(\frac{dr^2}{1-kr^2} + r^2(d\theta^2 + \sin^2 \theta d\phi^2) \right) \quad a(t): \text{scale factor}$$

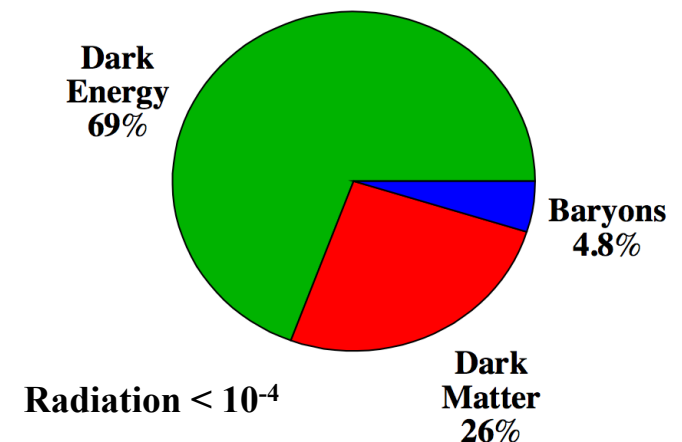
Friedmann equation :

$$\left(\frac{\dot{a}}{a} \right)^2 \equiv H^2 = \frac{8\pi G}{3} (\rho_R + \rho_M + \rho_\Lambda) - \frac{k}{a^2}$$

EoS: $p = \text{pressure} \equiv w \times \rho \Rightarrow \boxed{\rho \propto a^{-3(1+w)}}$

$$w = \begin{cases} 0 \text{ (matter)} & \Rightarrow a^{-3} \\ 1/3 \text{ (radiation)} & \Rightarrow a^{-4} \\ -1 \text{ } (\Lambda, \text{ dark energy}) & \Rightarrow a^0 \\ \text{Curvature} & \Rightarrow a^{-2} \end{cases}$$

Now.... ($a \equiv 1$)



Dynamics of the expanding Universe

Einstein equation & Friedmann-Lemaître-Robertson-Walker metrics

$$ds^2 = g_{\alpha\beta} dx^\alpha dx^\beta = dt^2 - a^2(t) \left(\frac{dr^2}{1-kr^2} + r^2(d\theta^2 + \sin^2 \theta d\phi^2) \right) \quad a(t): \text{scale factor}$$

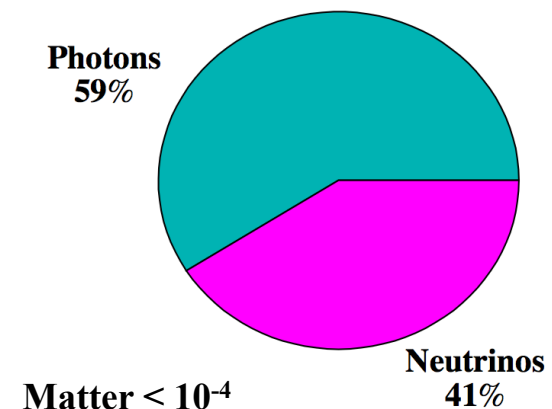
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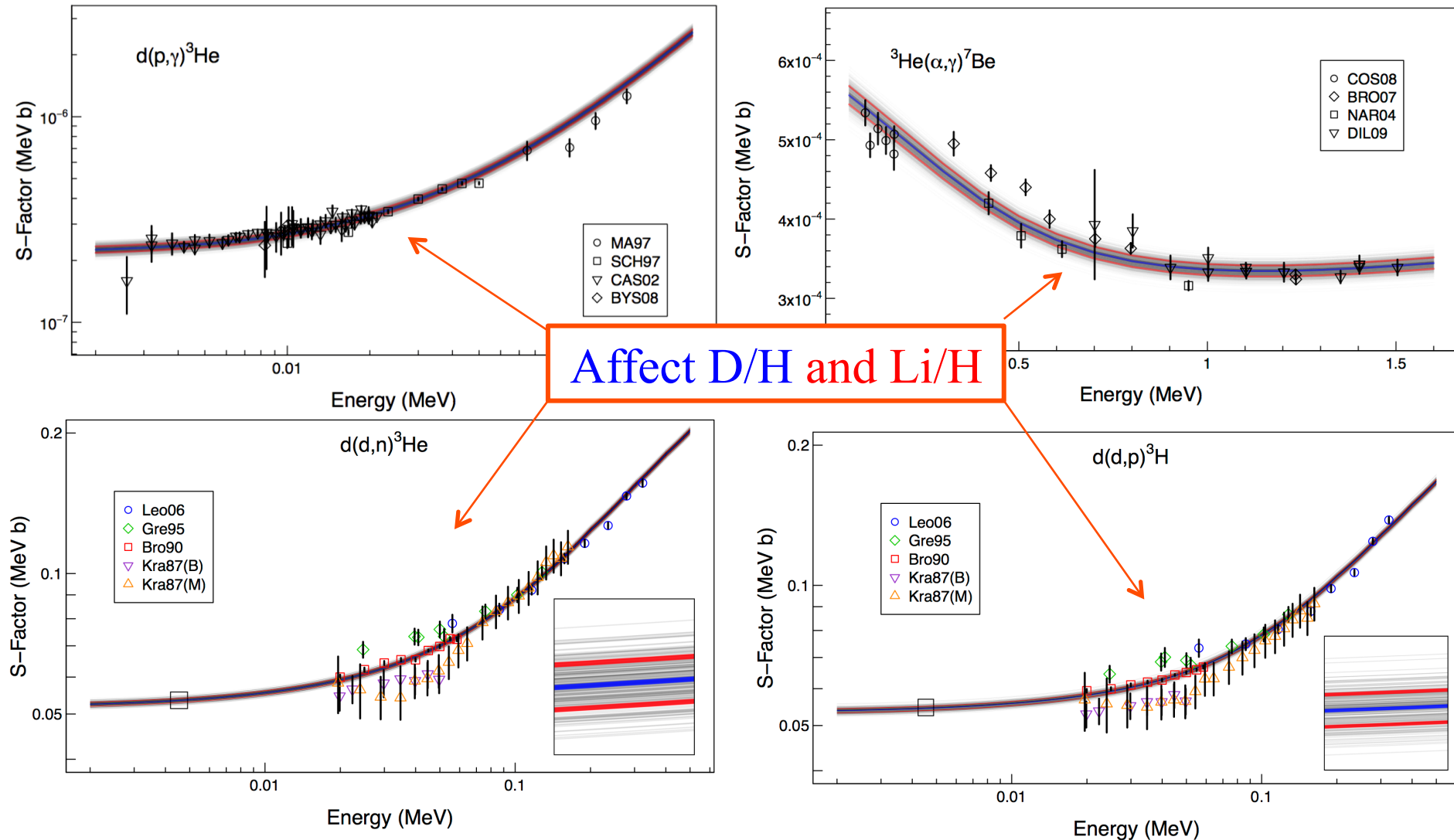
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....and then ($a \approx 10^{-8}$)



Re-evaluation of main reaction rates

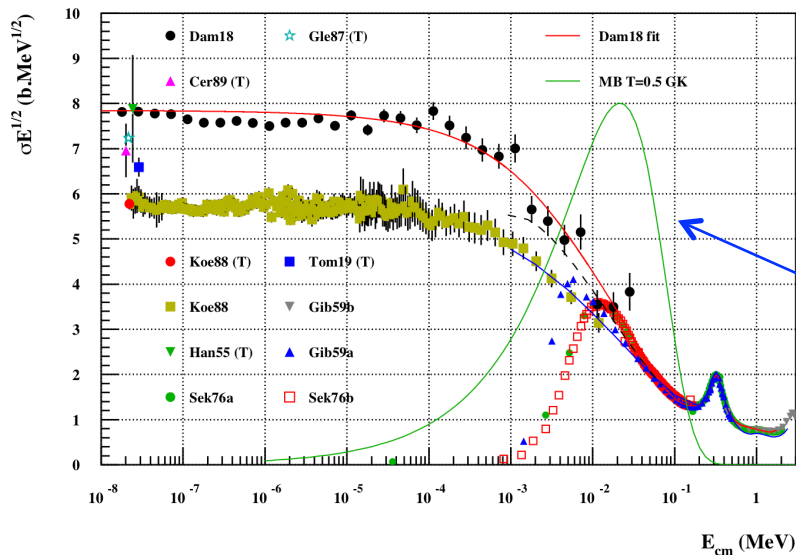
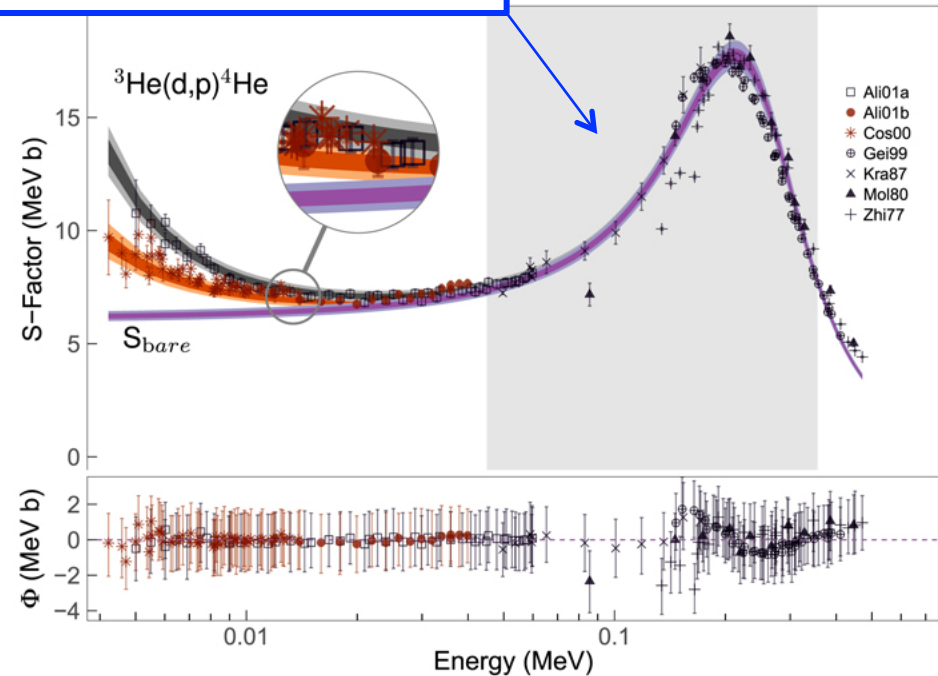
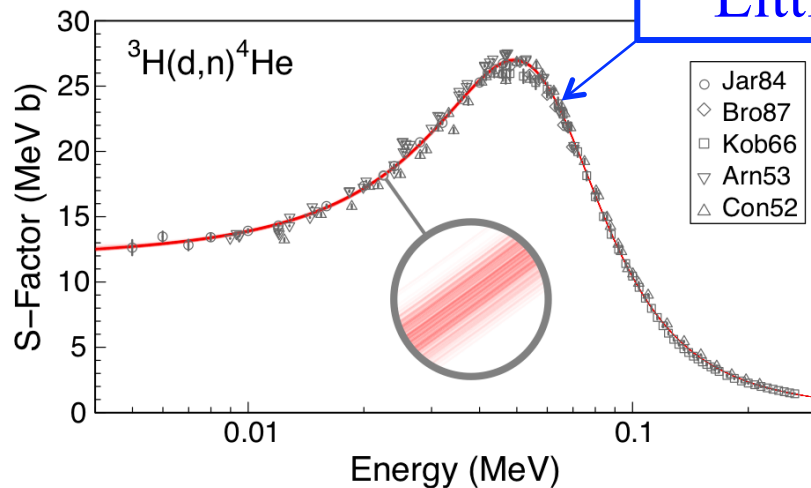
Bayesian analyses of reaction rates for BBN : $D(p,\gamma)^3\text{He}$, $D(d,n)^3\text{He}$, $D(d,p)^3\text{H}$ and $^3\text{He}(\alpha,\gamma)^7\text{Be}$ [*Iliadis+ 2016; Gómez Iñesta+2017*] already included in PRIMAT



Re-evaluation of main reaction rates

Bayesian analyses of reaction rates for BBN] including $^3\text{He}(d,p)^4\text{He}$, $^3\text{H}(d,n)^4\text{He}$ [de Souza+2019a,b] and $^7\text{Be}(n,p)^7\text{Li}$ [Tan Hong Kiat+ in preparation] to be included in PRIMAT

Little influence on BBN



Re-evaluation in progress
Direct influence on Li/H ($\lesssim 8\%$)

n $\leftrightarrow$ p weak reaction rates

$$n \rightarrow p : n \rightarrow p + e^- + \bar{\nu}_e \quad n + e^+ \rightarrow p + \bar{\nu}_e \quad n + \nu_e \rightarrow p + e^-$$

$$p \rightarrow n : p + e^- + \bar{\nu}_e \rightarrow n \quad p + \bar{\nu}_e \rightarrow n + e^+ \quad p + e^- \rightarrow n + \nu_e$$

$$\lambda_{n\leftrightarrow p} = C \sum \int (\text{phase space}) \times (\text{e distribution}) \times (\nu_e \text{ distribution}) dE$$

+ “some small corrections”

$$\lambda_{n\rightarrow pev} = C \int_1^q \frac{\varepsilon(\varepsilon - q)^2 (\varepsilon^2 - 1)^{1/2} d\varepsilon}{[1 + \exp(-\varepsilon z)] \{1 + \exp[(\varepsilon - q)z_\nu]\}} \quad T \rightarrow 0 \quad \boxed{\frac{1}{\tau_n} = C \int_1^q \varepsilon(\varepsilon - q)^2 (\varepsilon^2 - 1)^{1/2} d\varepsilon}$$

$$(q \equiv Q_{np}/m_e, \varepsilon \equiv E_e/m_e, z \equiv m_e/T_\gamma, z_\nu \equiv m_e/T_\nu)$$

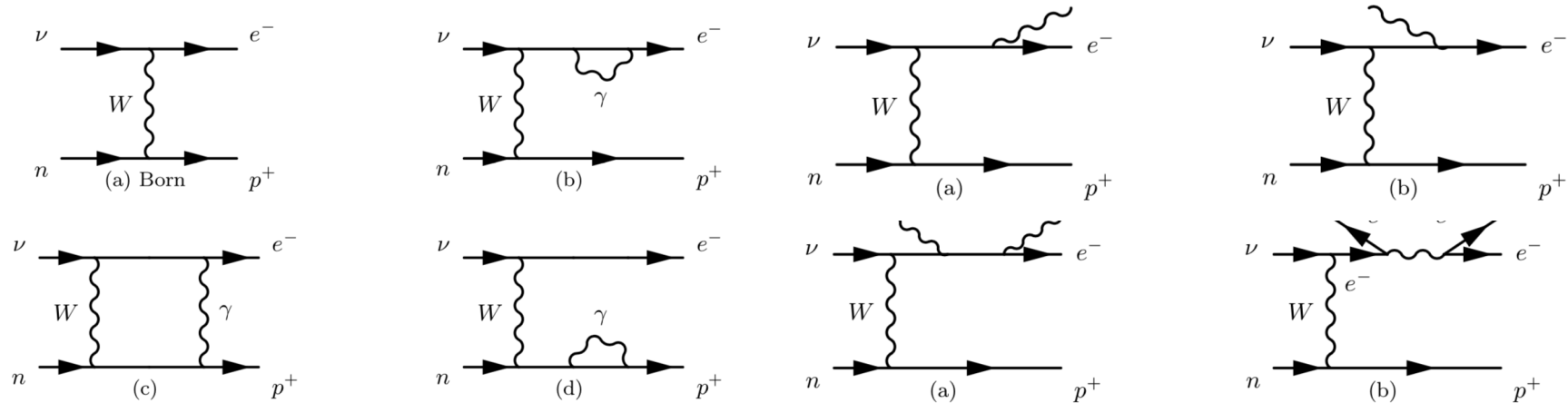
➤ Experimental neutron lifetime?

$$\Delta Y_p = +0.0002 \times \Delta \tau_n \text{ (s)}$$

➤ Calculation of the “small corrections”

“Small corrections” to the weak rates

1. radiative corrections
2. finite nucleon mass corrections,
3. finite temperature radiative corrections
4. weak-magnetism
5. QED plasma effects
6. incomplete neutrino decoupling



[Dicus+ 1982; Seckel 1993; Dolgov+ 1997; Lopez+ 1997; Lopez & Turner 1999; Brown & Sawyer 2001; Mangano+ 2005; Pisanti+ 2008; Grohs+ 2016; and many more]

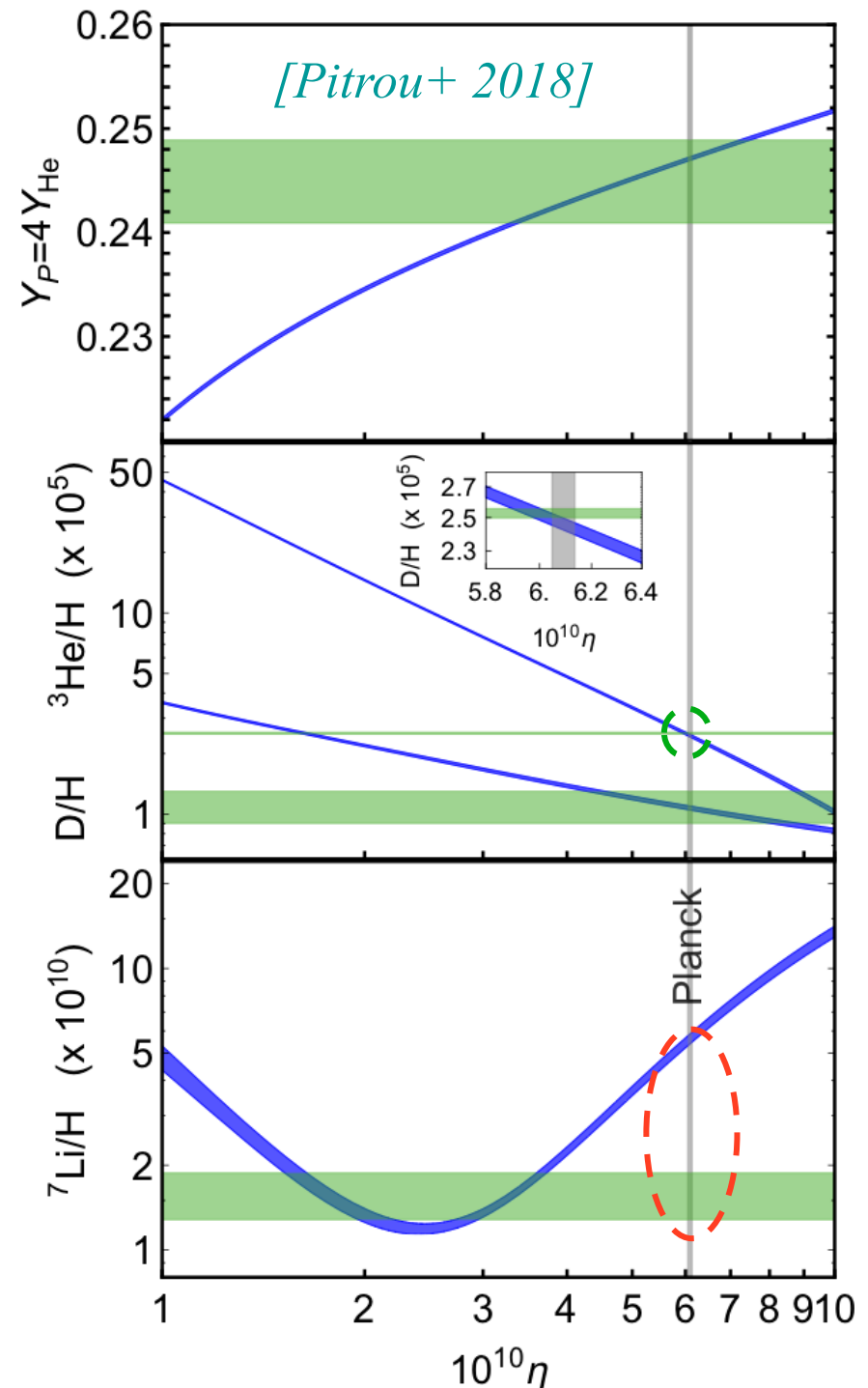
Pitrou et al. 2018 : All included and calculated in a self consistent way, allowing to take into account the correlations between them, and verifying that all satisfy detailed balance

Comparison between observed and calculated abundances

Limits ($1-\sigma$) obtained by Monte-Carlo fusing *Descouvemont+ 2004; Ando+ 2006, Iliadis+; 2016; Gómez Iñesta+ 2017;.....* reaction rate uncertainties.

Concordance (?) BBN, spectroscopy and CMB

- $\Omega_B h^2$ [*Planck: Ade+ 2016*]
- ^4He [*Aver+ 2015*]
- D [*Cooke+ 2018*]
- ^3He [*Bania et al. (2002)*]
- ^7Li [*Sbordone+ 2010*] : difference of a factor of ≈ 3 between calculated (BBN+CMB) and observed (Spite plateau) primordial lithium



Comparison between BBN codes

	BBN calculations			
	^4He	D/H	$^3\text{He}/\text{H}$	$^7\text{Li}/\text{H}$
	$\times 10^0$	$\times 10^{-5}$	$\times 10^{-5}$	$\times 10^{-10}$
Observations	0.2449 ± 0.0040	2.527 ± 0.030	$< (0.9-1.3)$	1.58 ± 0.31
<i>EZ_BBN (Coc+2015)</i>	0.2484 ± 0.0002	2.45 ± 0.05	1.07 ± 0.03	5.61 ± 0.26
<i>PRIMAT (Pitrou+2018)</i>	0.24709 ± 0.00017	2.459 ± 0.036	1.074 ± 0.026	5.623 ± 0.247

- Except for ^4He , very good agreement between (Fortran77) *EZ_BBN* and (Mathematica) *PRIMAT* results
- Small corrections on weak rates $\approx$ observational uncertainties!

	$\Delta^4\text{He}[\%]$	$\Delta(\text{D}/\text{H}) [\%]$
Radiative corrections	0.53	0.70
Finite mass of nucleons	1.27	0.25
Incomplete neutrino decoupling		0.37
QED effects on plasma		0.12
Total	1.84	1.49

Comparison between BBN codes

	BBN calculations			
	^4He	D/H	$^3\text{He}/\text{H}$	$^7\text{Li}/\text{H}$
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<i>Cyburt+2016</i>	0.24709 $\pm$ 0.00025	2.58 $\pm$ 0.13	1.0039 $\pm$ 0.0090	4.68 $\pm$ 0.67

Good agreement between Paris/Orsay and US (Cyburt) results

Comparison between BBN codes

	BBN calculations			
	^4He	D/H	$^3\text{He}/\text{H}$	$^7\text{Li}/\text{H}$
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<i>Cyburt+2016</i>	0.24709 ± 0.00025	2.58 ± 0.13	1.0039 ± 0.0090	4.68 ± 0.67
<i>Yeh priv. comm</i>	0.2472	2.449	1.076	5.633

Even better if one uses the same reaction rates!

(*Tsung-Han*) *Yeh priv. comm* = *Cyburt+2016* code with re-evaluated rates :

$\text{D}(p,\gamma)^3\text{He}$, $\text{D}(d,n)^3\text{He}$, $\text{D}(d,p)^3\text{H}$, $^7\text{Be}(n,\alpha)^4\text{He}$ & $^3\text{He}(\alpha,\gamma)^7\text{Be}$

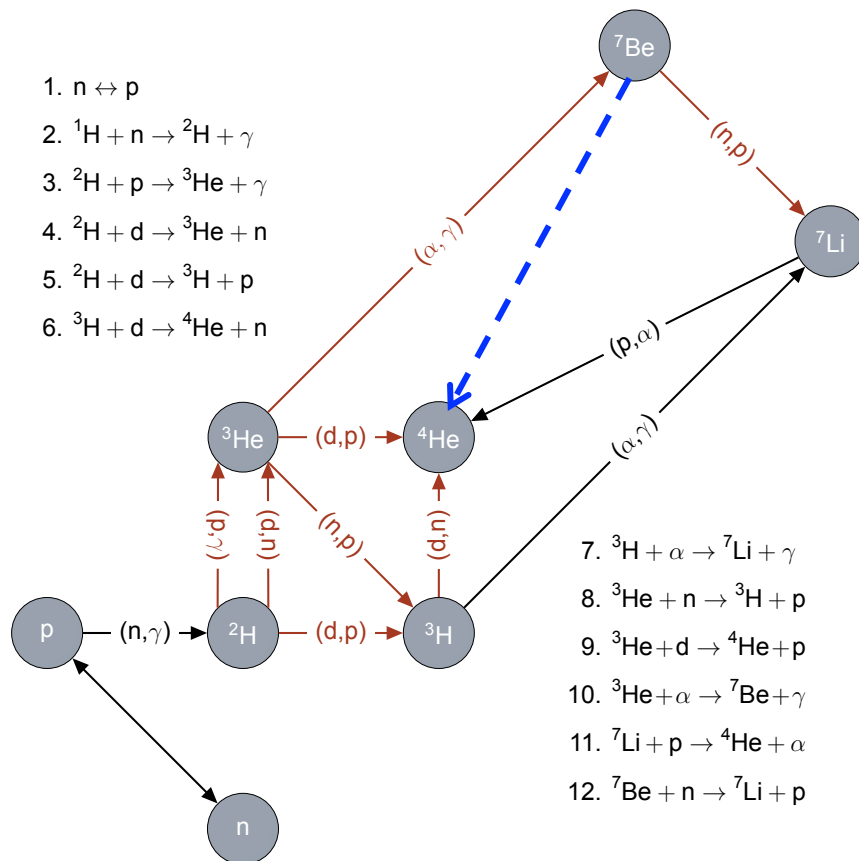
- The “**Lithium problem**”: difference of a factor of ≈ 3 between calculated and observed primordial lithium
- Precision needed for deuterium predictions (1% on reaction rates!)

Nuclear solution to the Li problem ?

At η_{CMB} ${}^7\text{Li}$ from ${}^7\text{Be}$ post BBN decay

Tentatives nuclear solutions:
 ${}^7\text{Be}$ destruction by:

1. Supplementary reactions
 e.g. ${}^7\text{Be}(d,p){}^8\text{Be}^* \rightarrow 2\alpha$



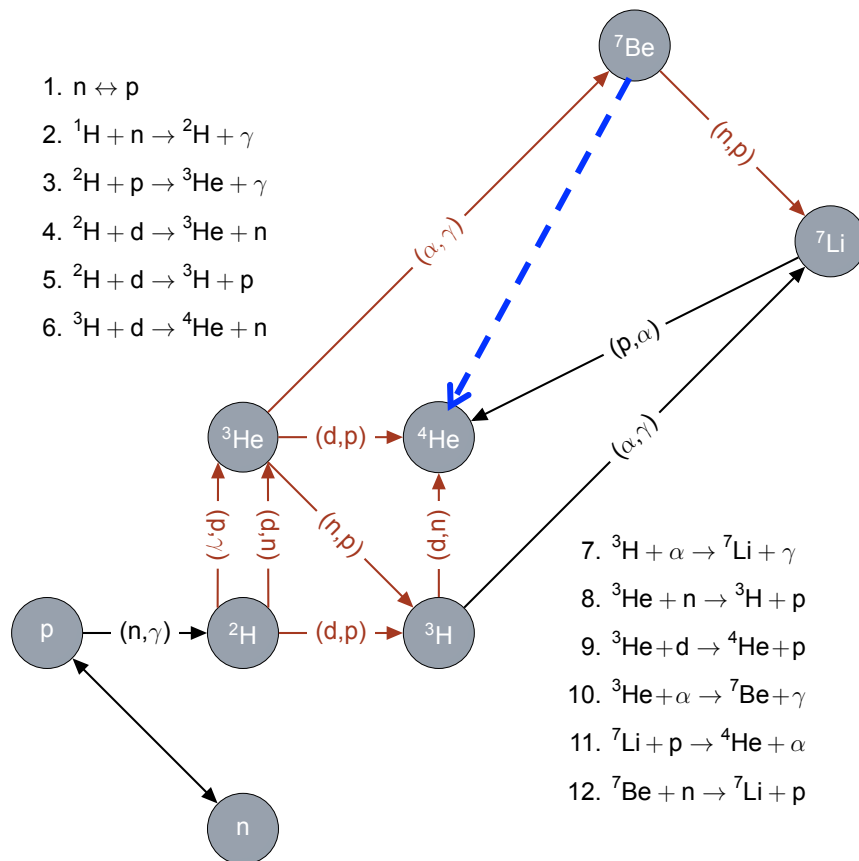
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Ruled out following extensive
 experimental and theoretical
 searches [Coc+ 2004; 2011, Angulo+
 2005, Kirsebom & Davids 2011,.....
 Rijal+ 2019].

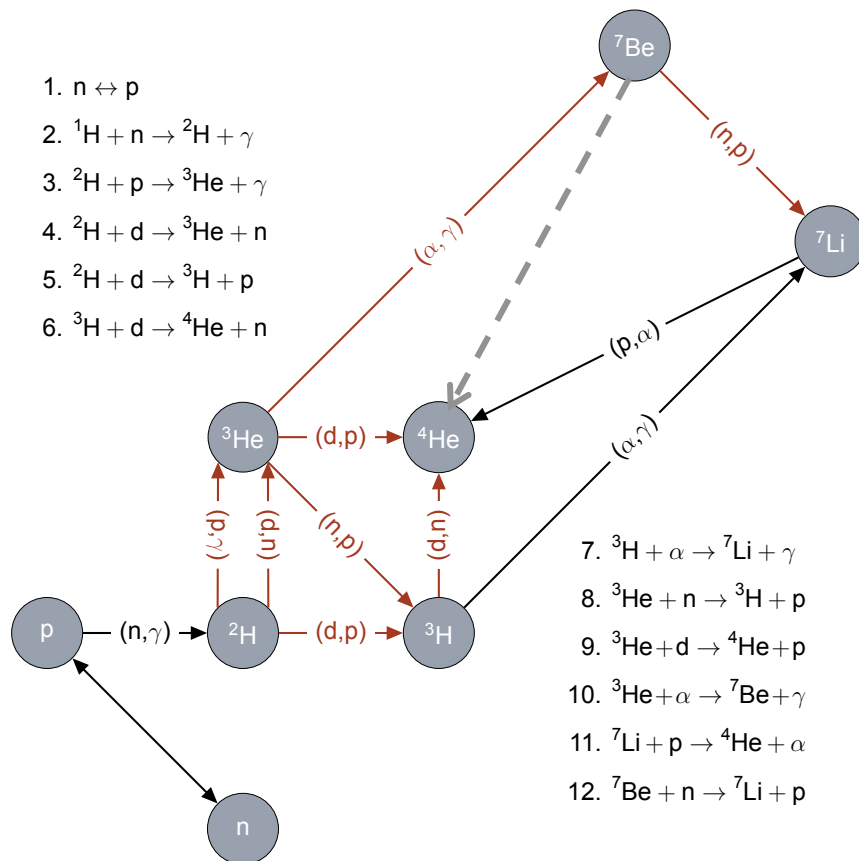


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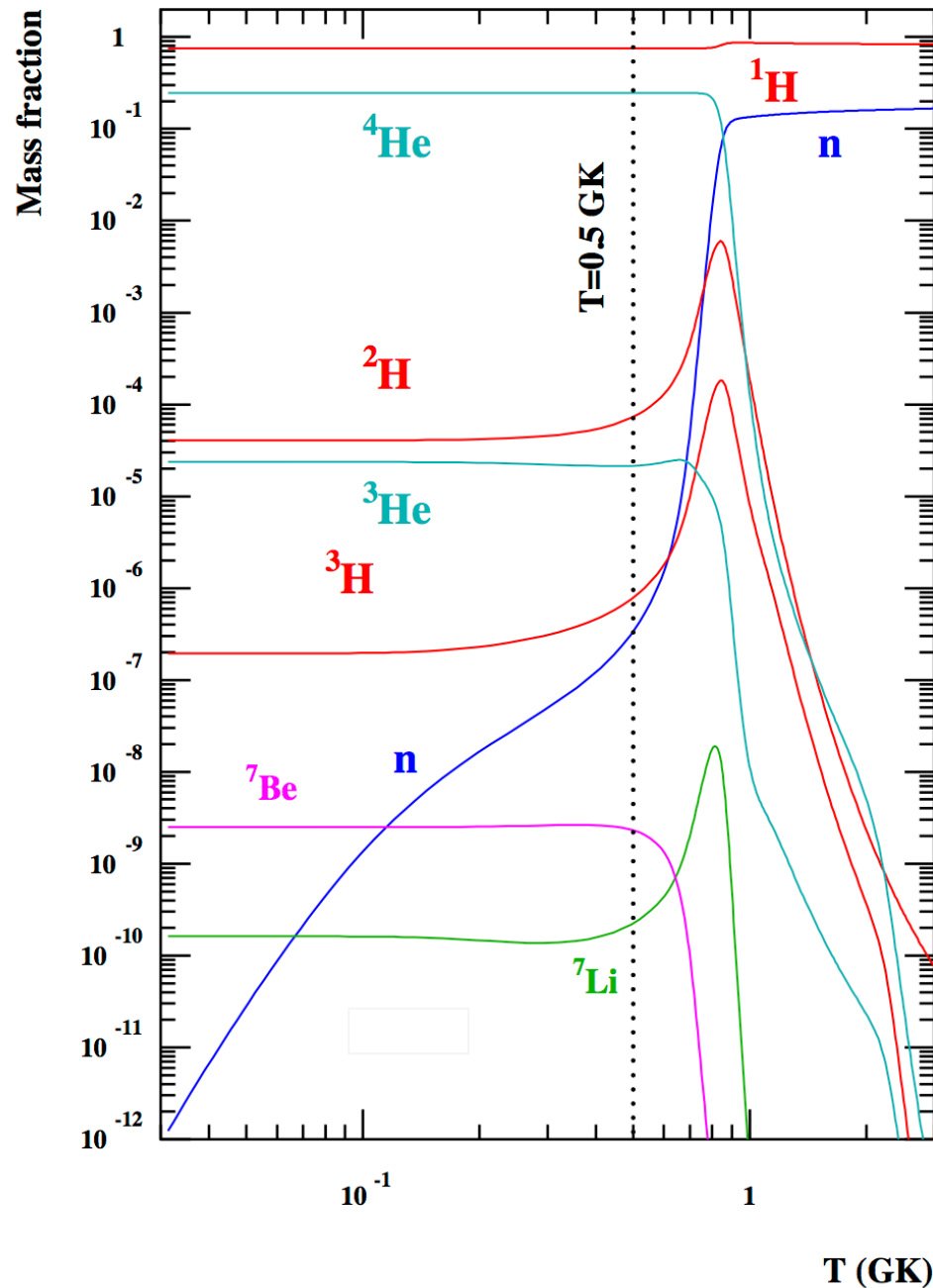
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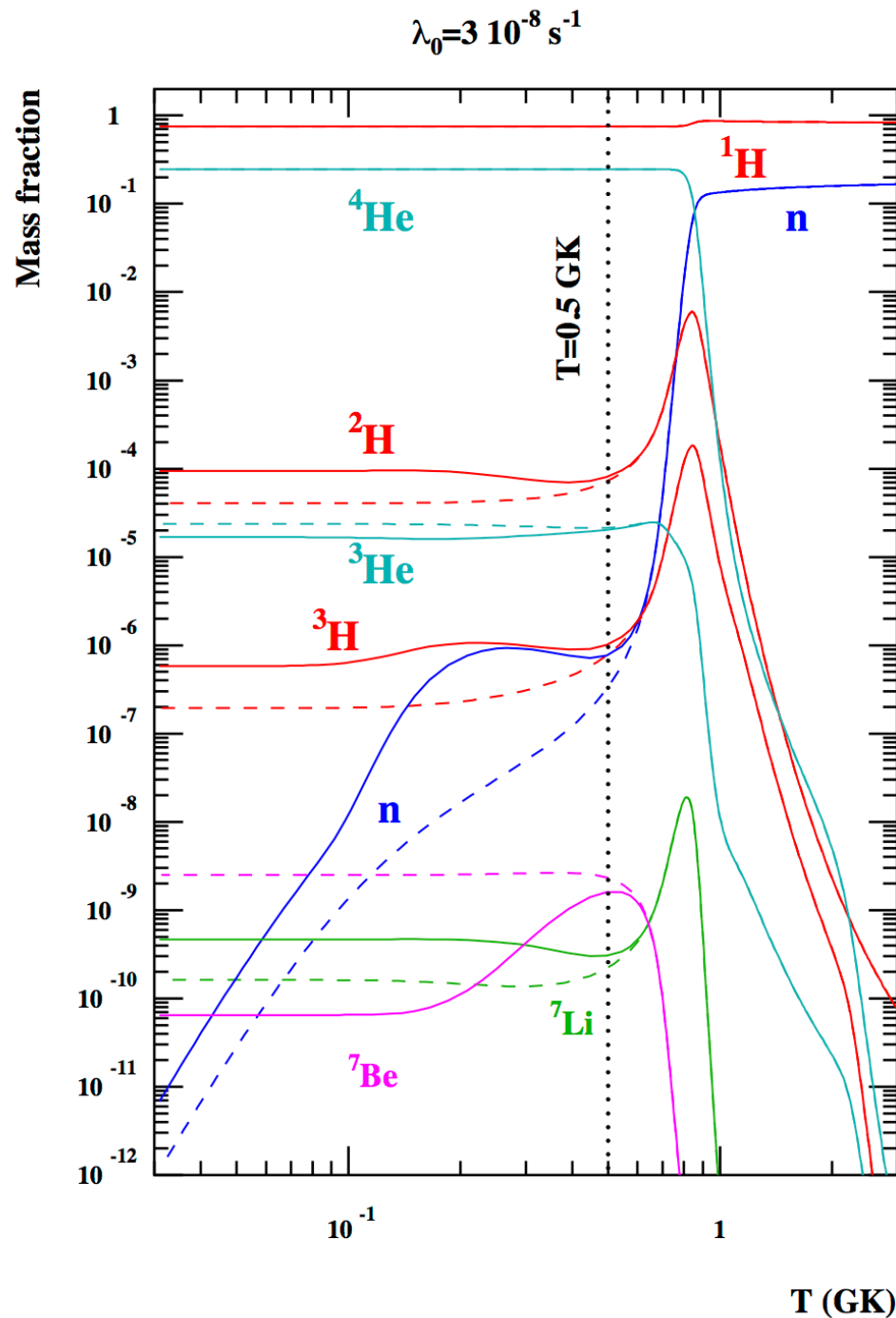
1. ~~Supplementary reactions
e.g. ${}^7\text{Be}(d,p){}^8\text{Be}^* \rightarrow 2\alpha$~~
2. Increased neutron destruction efficiency by ${}^7\text{Be}(\mathbf{n},p){}^7\text{Li}(p,\alpha){}^4\text{He}$ from exotic neutron sources



${}^7\text{Be}$ (${}^7\text{Li}$) destruction



- ☐ By most abundant projectiles
 ${}^7\text{Be}(p,\gamma){}^8\text{B}$ hindered by photo-dissociation ($Q=1.375$ MeV)
 and ${}^7\text{Be}(\alpha,\gamma){}^{11}\text{C}$ by Coulomb
- ☐ By less abundant projectiles
 but higher cross sections of
 overlooked reactions?
- ☐ No much activity below 0.5 GK
 (low charged particle rates and
 low neutron abundance).



Extra neutrons?

□ Increased neutron abundance at T above 1 GK only leads to higher ^4He ($p \rightarrow d \rightarrow t, ^3\text{He} \rightarrow ^4\text{He}$)

□ At lower T (< 0.5 GK)

➤ $^7\text{Be} \searrow$ by $^7\text{Be}(n, p)^7\text{Li}$

➤ $\text{D} \nearrow$ by $^1\text{H}(n, \gamma)\text{D}$

➤ $^3\text{H} \nearrow$ by $^3\text{He}(n, p)^3\text{H}$

➤ $^7\text{Li} \nearrow$ by $^3\text{H}(\alpha, \gamma)^7\text{Li}$ and $^7\text{Be}(n, p)^7\text{Li}$ [when $^7\text{Li}(p, \alpha)$ ineffective at low T]

➤ Minimum in $^7\text{Li} + ^7\text{Be}$ abundance

The limits to ${}^7\text{Li}+{}^7\text{Be}$ destruction by extra neutrons

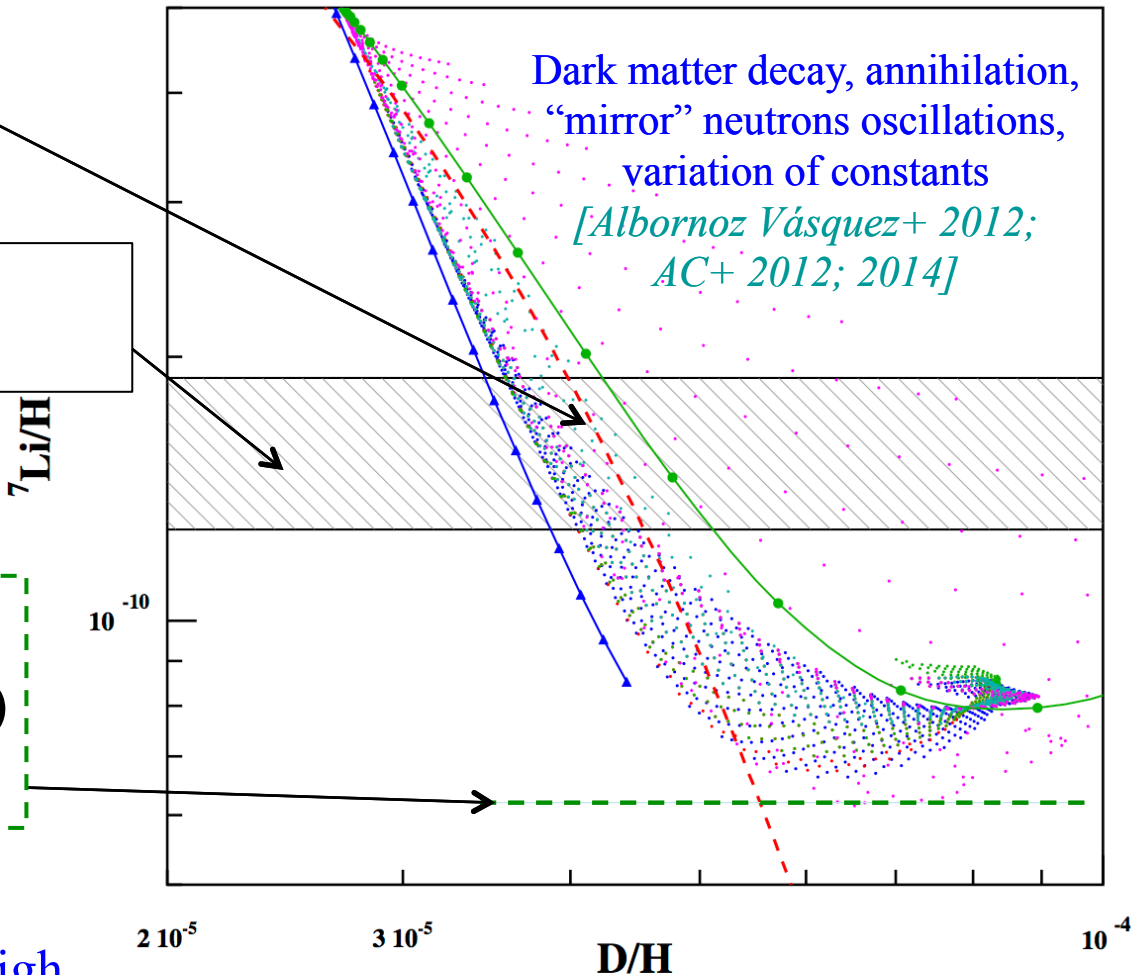
See also e.g. *Olive+ 2012; Kusakabe+ 2014;.....*



Li observational limits
[Sbordone+ 2010]

${}^3\text{He}(\text{n},\text{p})\,{}^3\text{H}(\alpha,\gamma){}^7\text{Li}$ and
 ${}^7\text{Be}(\text{n},\text{p}){}^7\text{Li}$ [when ${}^7\text{Li}(\text{p},\alpha)$
 ineffective at low T]

Dark matter decay, annihilation,
“mirror” neutrons oscillations,
variation of constants
[Albornoz Vázquez+ 2012;
AC+ 2012; 2014]



Even worse for (non-thermal) high energy neutrons [*Kusakabe+ 2004*]

Lower Li/H $\Rightarrow$ higher D/H

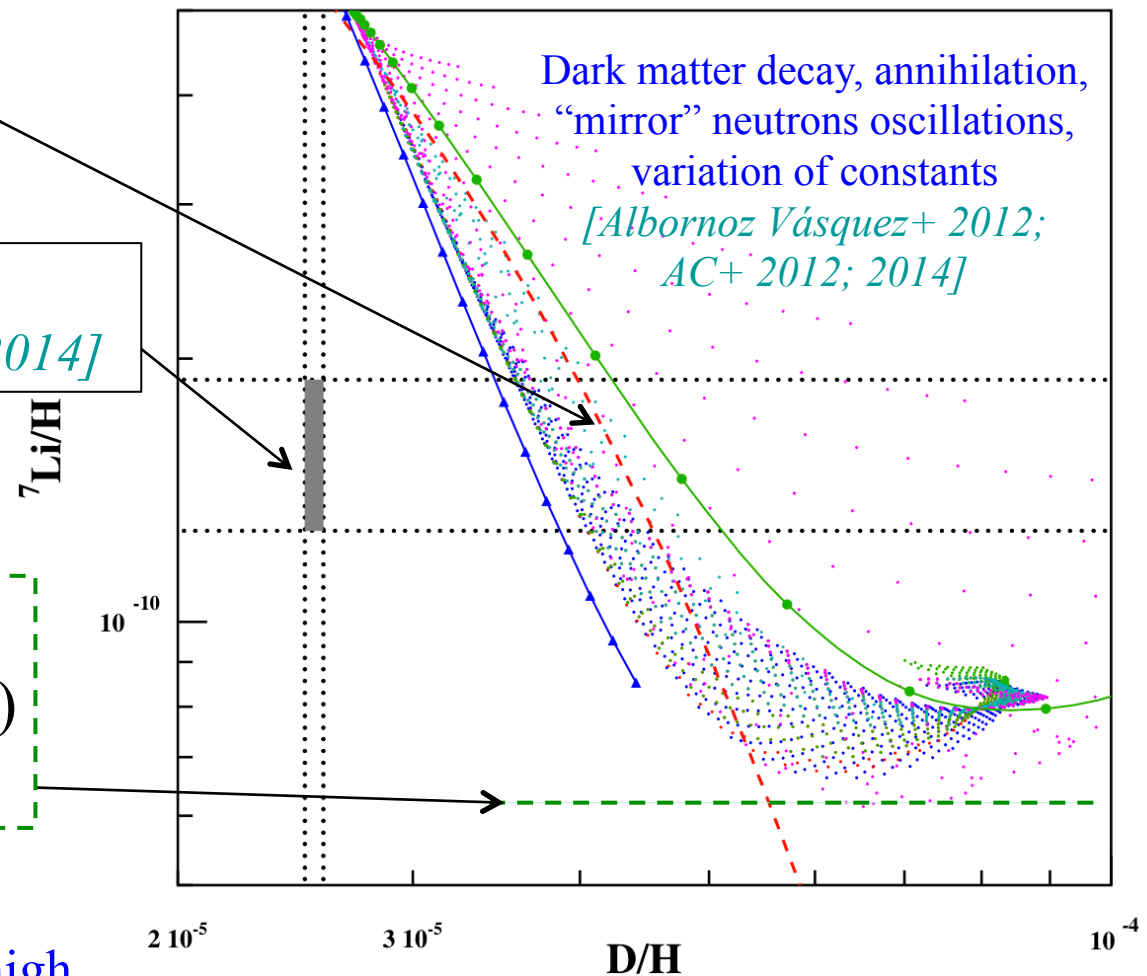
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See also e.g. *Olive+ 2012; Kusakabe+ 2014;*



Li/D observational limits
[*Sbordone+ 2010* $\times$ *Cooke+ 2014*]

${}^3\text{He}(\text{n},\text{p}) {}^3\text{H}(\alpha,\gamma) {}^7\text{Li}$ and
 ${}^7\text{Be}(\text{n},\text{p}) {}^7\text{Li}$ [when ${}^7\text{Li}(\text{p},\alpha)$
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<http://www2.iap.fr/users/pitrou/primat.htm>

Precision big bang nucleosynthesis with improved Helium-4 predictions

CSNSM+IAP

Cyril Pitrou^{a,b,*}, Alain Coc^c, Jean-Philippe Uzan^{a,b}, Elisabeth Vangioni^{a,b}

^a Institut d'Astrophysique de Paris, CNRS UMR 7095, 98 bis Bd Arago, 75014 Paris, France

^b Sorbonne Université, Institut Lagrange de Paris, 98 bis Bd Arago, 75014 Paris, France

^c Centre de Sciences Nucléaires et de Sciences de la Matière (CSNSM), CNRS IN2P3, Univ. Paris-Sud, Université Paris-Saclay, Bâtiment 104, F-91405 Orsay Campus, France



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STANDARD BIG BANG NUCLEOSYNTHESIS UP TO CNO WITH AN IMPROVED EXTENDED NUCLEAR NETWORK

ALAIN COC¹, STÉPHANE GORIELY², YI XU², MATTHIAS SAIMPET³, AND ELISABETH VANGIONI³

¹ Centre de Spectrométrie Nucléaire et de Spectrométrie de Masse (CSNSM), CNRS/IN2P3, Université Paris Sud, UMR 8609, Bâtiment 104, F-91405 Orsay Campus, France

² Institut d'Astronomie et d'Astrophysique, Université Libre de Bruxelles, CP 226, Boulevard du Triomphe, B-1050 Bruxelles, Belgium

³ Institut d'Astrophysique de Paris, UMR 7095 CNRS, Université Pierre et Marie Curie, 98 bis Boulevard Arago, Paris 75014, France

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Compilation and *R*-matrix analysis of Big Bang nuclear reaction rates

Pierre Descouvemont^{a,*}, Abderrahim Adahchour^{a,2}, Carmen Angulo^b,
Alain Coc^c, Elisabeth Vangioni-Flam^d

^a Physique Nucléaire Théorique et Physique Mathématique, CP229, Université Libre de Bruxelles, B-1050 Brussels, Belgium

^b Centre de Recherches du Cyclotron, Université catholique de Louvain, Chemin du Cyclotron 2, B-1348 Louvain-la-Neuve, Belgium

^c Centre de Spectrométrie Nucléaire et de Spectrométrie de Masse, CNRS/IN2P3/UPS, Bât. 104, F-91405 Orsay Campus, France

^d Institut d'Astrophysique de Paris, CNRS, 98^{bis} Bd. Arago, 75014 Paris, France

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THE ${}^7\text{Be}(d,p){}^6\text{Li}$ CROSS SECTION AT BIG BANG ENERGIES AND THE PRIMORDIAL ${}^7\text{Li}$ ABUNDANCE

C. ANGULO, E. CASAREJOS, M. COUDER,¹ P. DEMARET, P. LELEUX,² AND F. VANDERRIST
Centre de Recherches du Cyclotron and Institut de Physique Nucléaire, Université catholique de Louvain, Chemin du Cyclotron 2,
B-1348 Louvain-la-Neuve, Belgium; angulo@cyc.ucl.ac.be

A. COC, J. KIENER, AND V. TATISCHEFF
Centre de Spectrométrie Nucléaire et de Spectrométrie de Masse, IN2P3-CNRS, Université Paris-Sud, Bâtiment 104, F-91405 Orsay Campus, France

T. DAVINSON AND A. S. MURPHY
School of Physics, James Clerk Maxwell Building, King's Buildings, University of Edinburgh, Mayfield Road, Edinburgh EH9 3JZ, UK

N. L. ACHOURI AND N. A. ORR
Laboratoire de Physique Corpusculaire, ENSICAEN, and Université de Caen, IN2P3-CNRS, 6 boulevard du Maréchal Juin,
F-14050 Caen Cedex, France

D. CORTINA-GIL
Departamento de Física de Partículas, Universidad de Santiago de Compostela, E-15782 Santiago de Compostela, Spain

P. FIGUERA
Laboratori Nazionali del Sud, Istituto Nazionale di Fisica Nucleare, via Santa Sofia 62, I-95123 Catania, Italy

B. R. FULTON
Department of Physics, University of York, Heslington, York YO10 5DD, UK

I. MUKHA³
Institut voor Kern- en Stralingsfysica, Katholieke Universiteit Leuven, Celestijnenlaan 200D,
B-3001 Leuven, Belgium

AND
E. VANGIONI
Institut d'Astrophysique de Paris, 98bis boulevard Arago, F-75014 Paris, France
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Conclusions

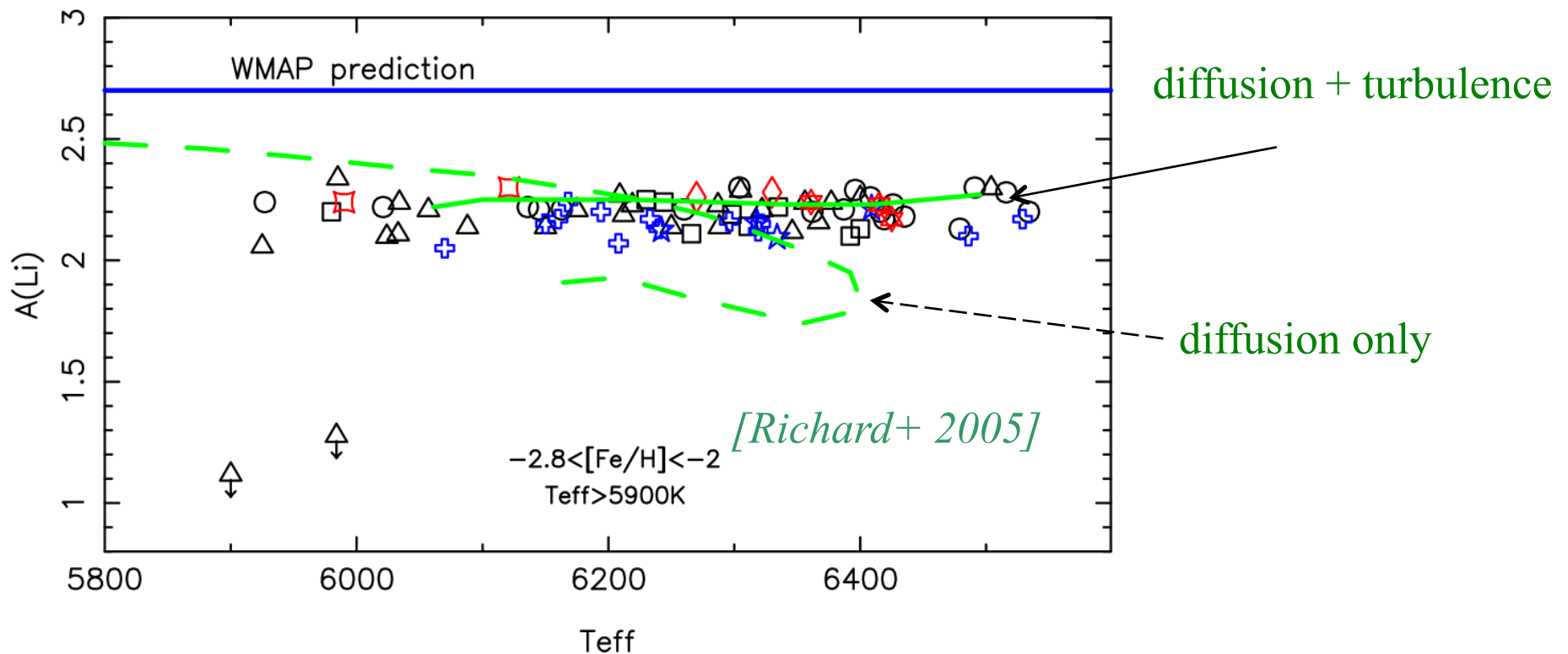
- ❑ Standard BBN is now in the (1%) precision era for D and ^4He
 - Precision deuterium observations (plateau ?) call for
 - even better precision on $\text{D}(\text{p},\gamma)^3\text{He}$, $\text{D}(\text{d},\text{n})^3\text{He}$ and $\text{D}(\text{d},\text{p})^3\text{H}$ cross sections
 - Corrections to the weak rates and improved neutron lifetime for ^4He
- ❑ However the lithium problem is worse than ever!
 - Disagreement (factor of 3) with Li observations
 - Nuclear : excluded by experiments
 - Cosmology or particle physics solutions overproduce deuterium
 - Stellar depletion, seemingly unavoidable, needs to be uniform
- ❑ Convergence of BBN codes when same nuclear reaction rates are used
 - Mathematica versus (independent) Fortran versions
 - Mathematica code with >400 reaction network publicly available at <http://www2.iap.fr/users/pitrou/primat.htm>

Li stellar depletion ?

□ In situ destruction (atomic diffusion + turbulence interplay)

➤ Some amount of depletion, a factor of 1.5 to 2, is unavoidable because of atomic diffusion [*Michaud+ 1984*]

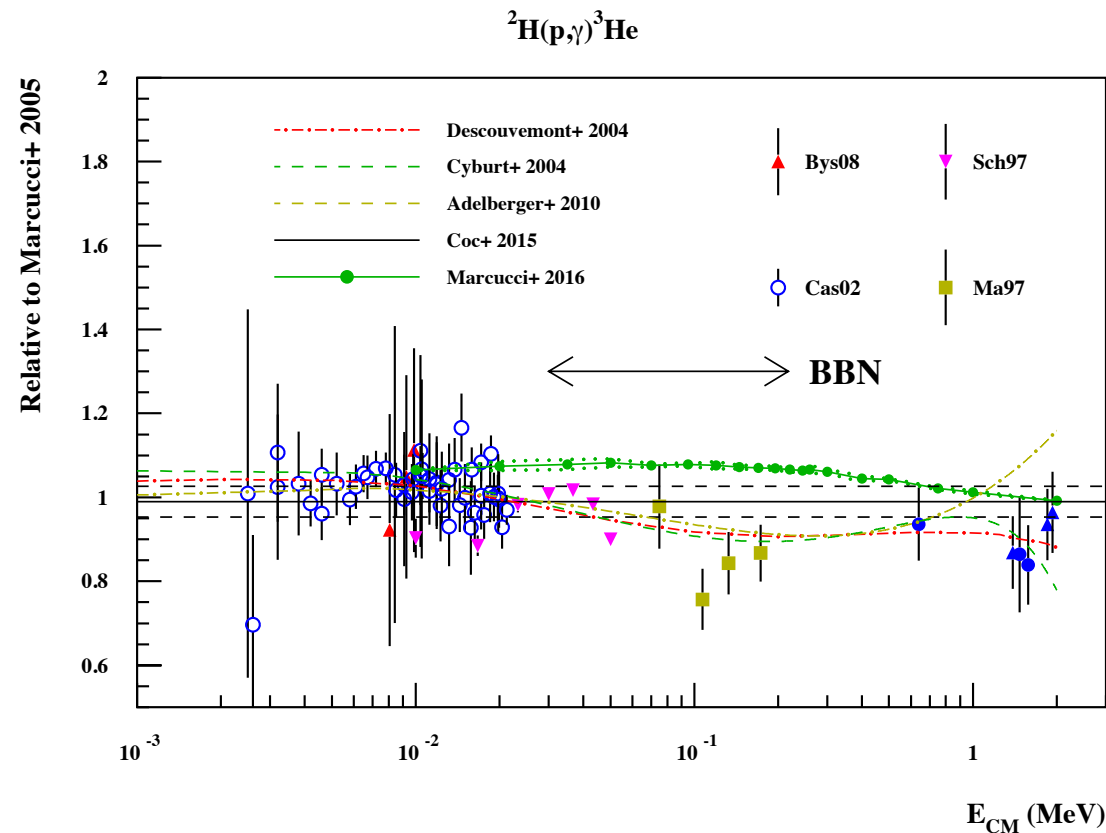
➤ Uniformity restored by an additional mixing process [*Richard+ 2005; Korn+ 2006*]



D/H sensitivity to reaction rates

$$\frac{\Delta(D/H)}{D/H} = -0.32 \times \frac{\Delta\langle\sigma v\rangle_{d(p,\gamma)^3\text{He}}}{\langle\sigma v\rangle_{d(p,\gamma)^3\text{He}}} - 0.54 \times \frac{\Delta\langle\sigma v\rangle_{d(d,n)^3\text{He}}}{\langle\sigma v\rangle_{d(d,n)^3\text{He}}} - 0.46 \times \frac{\Delta\langle\sigma v\rangle_{d(d,p)^3\text{H}}}{\langle\sigma v\rangle_{d(d,p)^3\text{H}}}$$

$D(p,\gamma)^3\text{He}$, $D(d,n)^3\text{He}$ and $D(d,p)^3\text{H}$ reaction rates need to be known at a few % level to match the 1.6% precision on observations!



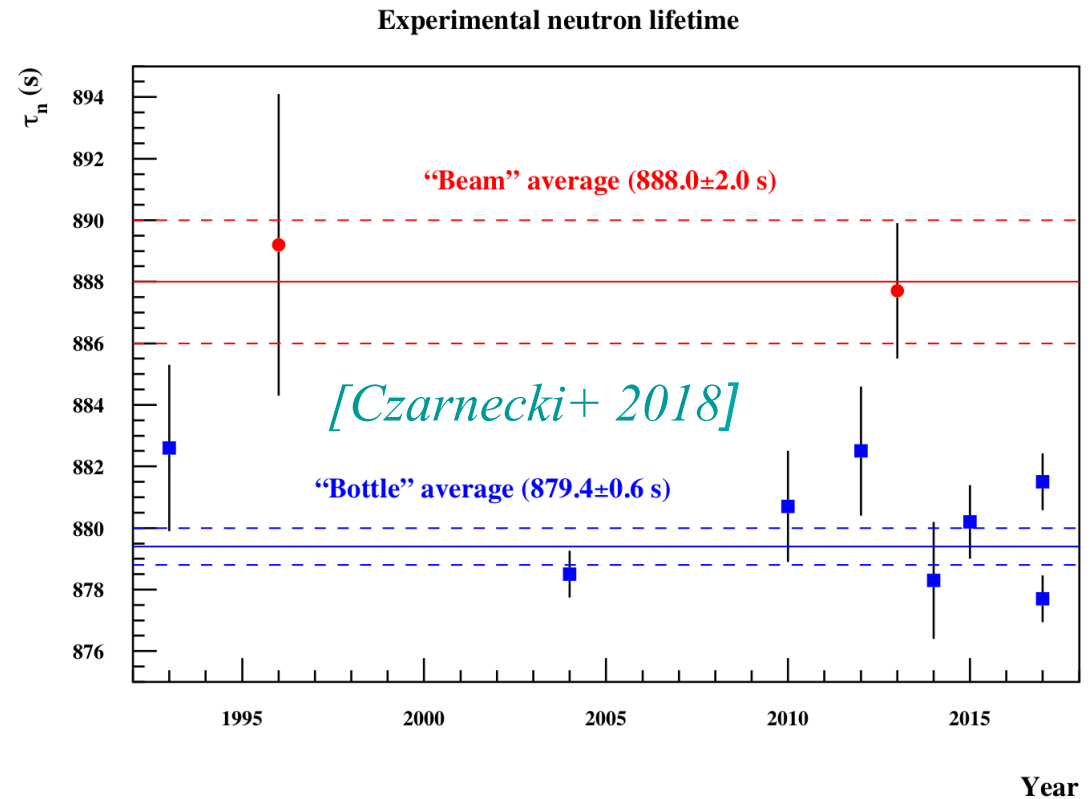
Neutron lifetime and the weak rates

Significant (3.8σ [Pignol 2015]) difference between:

- “beam” (protons from **beta decay only** counting) and
- “bottle” (surviving neutron counting)

experiments originating from

- unknown systematics or
- neutron disappearance in an other (brane/mirror) world [Sarazin+ 2012] or decay to DM [Fornal & Grinstein 2018]



In BBN governed by “surviving” neutrons like in “bottle” experiments

Using $\tau_n = 879.5 \pm 0.8$ s [Serebrov+ 2017; W. Marciano priv. comm.], close to average of “bottle” experiments ($\tau_n = 880.2 \pm 1.0$ s in PDG 2017)

