

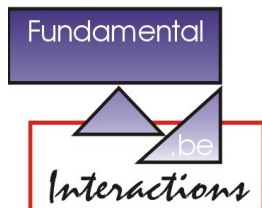
Neutrino magnetic moments

Julian Heeck

Solvay Workshop

'Beyond the Standard Model with Neutrinos and Nuclear Physics'

30.11.2017



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ULB

Why?

Abschrift

Physikalisches Institut
der Eidg. Technischen Hochschule
Zürich

Zürich, 4. Dez. 1930
Gloriastrasse

Liebe Radioaktive Damen und Herren,

Wie der Ueberbringer dieser Zeilen, den ich huldvollst
ansuhören bitte, Ihnen des näheren auseinandersetzen wird, bin ich

...

Nun handelt es sich weiter darum, welche Kräfte auf die
Neutronen wirken. Das wahrscheinlichste Modell für das Neutron scheint
mir aus wellenmechanischen Gründen (näheres weiss der Ueberbringer
dieser Zeilen) dieses zu sein, dass das ruhende Neutron ein
magnetischer Dipol von einem gewissen Moment μ ist. Die Experimente
verlangen wohl, dass die ionisierende Wirkung eines solchen Neutrons
nicht grösser sein kann, als die eines gamma-Strahls und darf dann
 μ wohl nicht grösser sein als $e \cdot (10^{-13} \text{ cm})$.

...

ges. W. Pauli

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Now it is also a question of which forces act upon neutrons. For me, the most likely model for the neutron seems to be, for wave-mechanical reasons (the bearer of these lines knows more), that the neutron at rest is a **magnetic dipole with a certain moment μ** . The experiments seem to require that the ionizing effect of such a neutron can not be bigger than the one of a gamma-ray, and then μ is probably not allowed to be larger than $e \cdot (10^{-13} \text{cm})$.

Nun handelt es sich weiter darum, welche Kräfte auf die Neutronen wirken. Das wahrscheinlichste Modell für das Neutron scheint mir aus wellenmechanischen Gründen (näheres weiss der Ueberbringer dieser Zeilen) dieses zu sein, dass das ruhende Neutron ein magnetischer Dipol von einem gewissen Moment μ ist. Die Experimente verlangen wohl, dass die ionisierende Wirkung eines solchen Neutrons nicht grösser sein kann, als die eines gamma-Strahls und darf dann μ wohl nicht grösser sein als $e \cdot (10^{-13} \text{cm})$.

$\sim 0.01 \mu_B$.

...

ges. W. Pauli

Why?

- Neutrino masses & mixing = solid evidence for BSM.
- Implies other observables, such as (for Dirac ν)

- Lepton flavor violation:

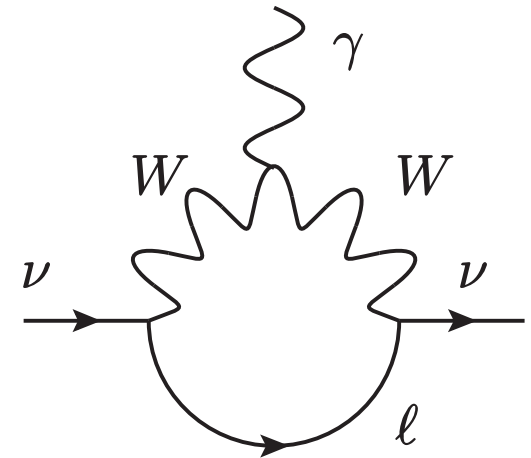
$$\frac{\Gamma(\ell_\alpha \rightarrow \ell_\beta \gamma)}{\Gamma(\ell_\alpha \rightarrow \ell_\beta \nu_\alpha \bar{\nu}_\beta)} \simeq \frac{3\alpha_{\text{EM}}}{32\pi} \left| \sum_{j=2,3} U_{\alpha j} \frac{\Delta m_{j1}^2}{M_W^2} U_{j\beta}^\dagger \right|^2 < 5 \times 10^{-53}.$$

[Petcov, '77; Cheng & Li, '77]

- Neutrino magnetic moment:

$$\mu_{ii}^D = \frac{3eG_F m_i}{8\sqrt{2}\pi^2} \simeq 3 \times 10^{-19} \mu_B \left(\frac{m_i}{\text{eV}} \right).$$

[Fujikawa, Shrock, '80]

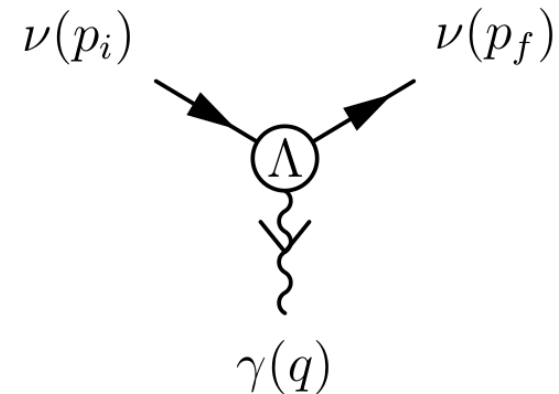


Observation = physics beyond m_ν !

Framework

- General interaction of ν mass eigenstates with photon A :

$$\mathcal{H} = \sum_{i,j=1}^3 \bar{\nu}_f \Lambda_{\mu}^{fi} \nu_i A^{\mu}.$$



$$(\gamma_{\mu} - q_{\mu} \not{q} / q^2) [f_Q(q^2) + f_A(q^2) q^2 \gamma_5] - i \sigma_{\mu\nu} q^{\nu} [f_M(q^2) + i f_E(q^2) \gamma_5]$$

charge

anapole

magnetic

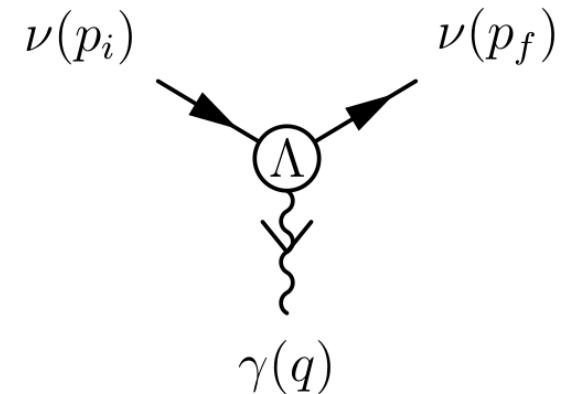
electric

- Hermitian form factor matrices f_x give **moments** $f_x(0)$:

$$\mu = f_M(0), \epsilon = f_E(0), \dots$$

- For Majorana $\nu = \nu^c$: $f_A = +f_A^T$, $f_{Q,M,E} = -f_{Q,M,E}^T$.

Framework



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- Hermitian form factor matrix

$$\mu = f_M(0)$$

$$\mu_{\text{Maj.}} = \begin{pmatrix} 0 & \mu_{e\mu} & \mu_{e\tau} \\ -\mu_{e\mu} & 0 & \mu_{\mu\tau} \\ -\mu_{e\tau} & -\mu_{\mu\tau} & 0 \end{pmatrix}.$$

- For Majorana $\nu = \nu^c$: $f_A = +f_A^T$, $f_{Q,M,E} = -f_{Q,M,E}^T$.

μ_ν in cosmology and astrophysics

- γ has plasma mass $\propto \sqrt{n_e/m_e}$.

\Rightarrow Plasmon decay: $\gamma^* \rightarrow \nu\bar{\nu}$.

- New **cooling** channel for stars!

[Bernstein+, '63; Raffelt, '90s; Viaux+, '13]

$$\sqrt{\sum_{i,j} |\mu_{ij}|^2} < 4.5 \times 10^{-12} \mu_B.$$

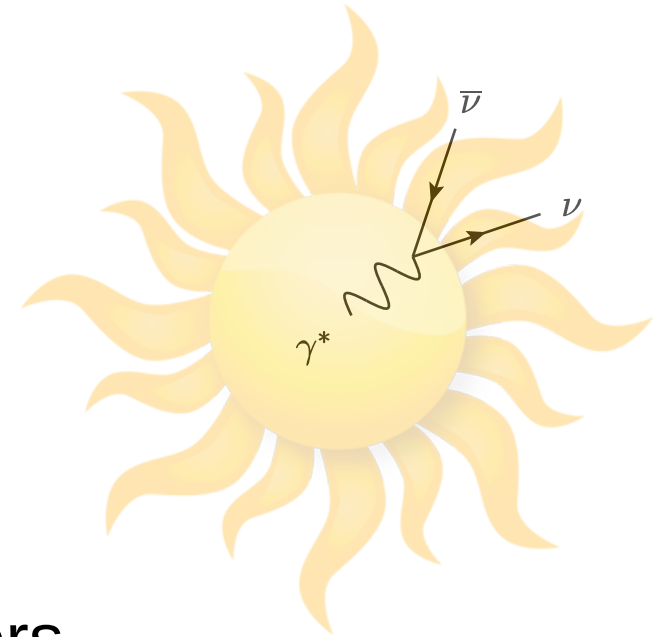
from red-giant branch in globular clusters.

- Big Bang Nucleosynthesis: $|\mu_{ij}| < 6 \times 10^{-10} \mu_B$.

[Vassh+ '15]

- For Dirac: $e \nu_L \rightarrow e \nu_R$ (in SN1987): $|\mu_{ij}| \lesssim 5 \times 10^{-13} \mu_B$.

[Morgan, '81; Fukugita, Yazaki, '87; Barbieri, Mohapatra, '88; Ayala+, '99; Kuznetsov, Mikheev, '07]

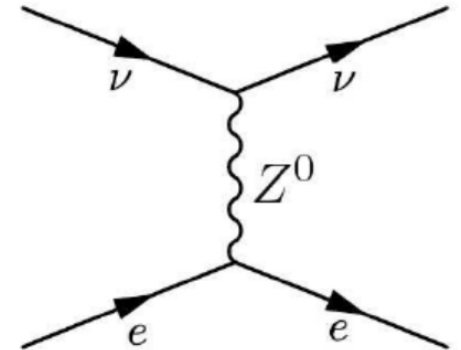


Magnetic moment in lab

- Clean probe: elastic $\nu_\alpha + e \rightarrow \nu_\beta + e$.

- Observable recoil energy T_e .

- Incoherent: $\frac{d\sigma}{dT_e} = \left(\frac{d\sigma}{dT_e}\right)_{SM} + \left(\frac{d\sigma}{dT_e}\right)_{\mu\nu}$.



$$r_1 + r_2 T_e + r_3 T_e^2$$

$$\frac{\pi\alpha^2}{m_e^2} \left(\frac{\mu_\nu}{\mu_B}\right)^2 \left(\frac{1}{T_e} - \frac{1}{E_\nu}\right)$$

- μ_ν wins for

[Engel, Vogel, '89]

$$T_e \lesssim \frac{\pi^2\alpha^2}{G_F^2 m_e^3} \left(\frac{\mu_\nu}{\mu_B}\right)^2$$

$$\simeq \text{few keV} \left(\frac{\mu_\nu}{10^{-11}\mu_B}\right)^2.$$

Need low thresholds!

Magnetic moment in lab

- Clean probe: elastic $\nu_e + e \rightarrow \nu_e + e$.

- Observable recoil

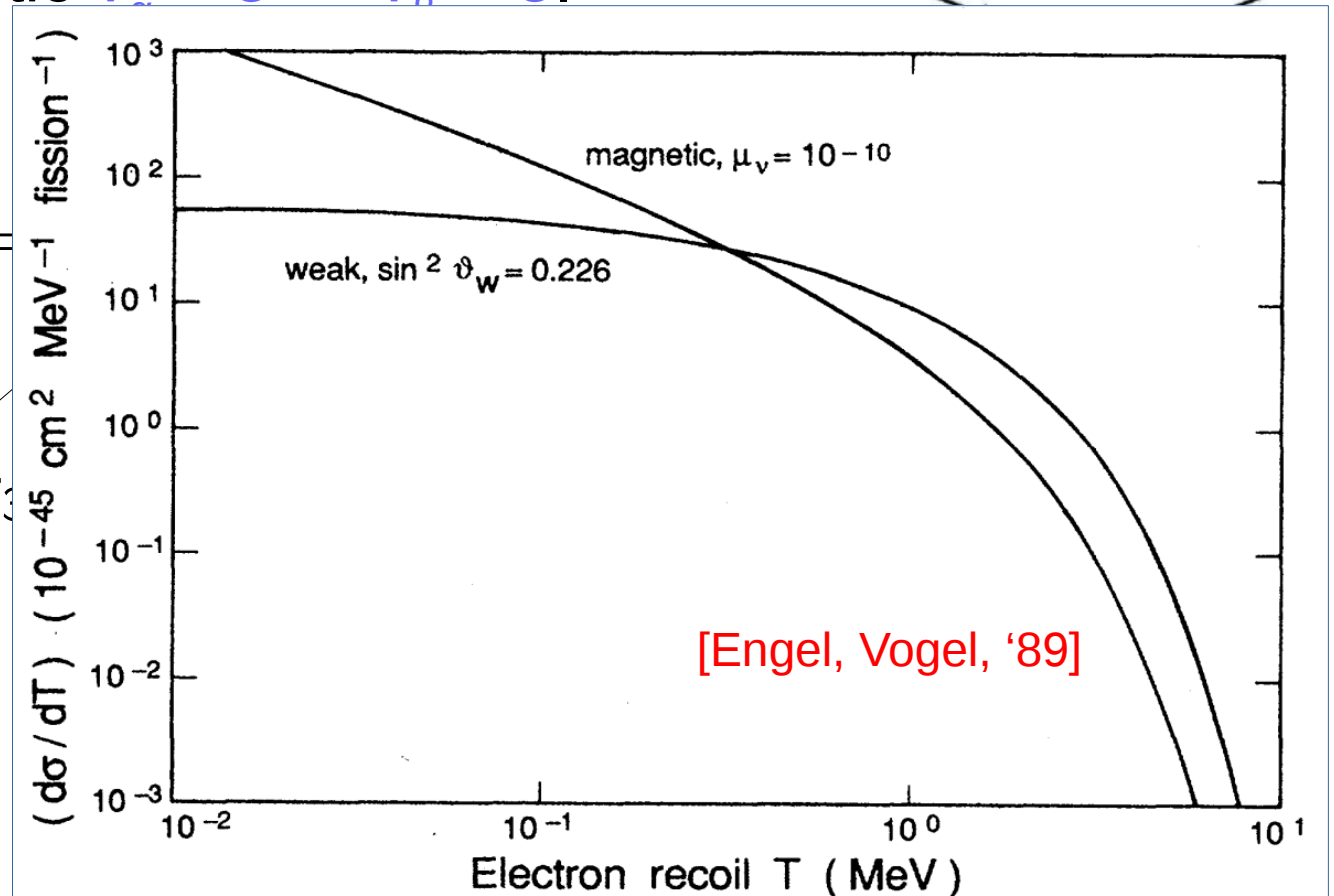
- Incoherent: $\frac{d\sigma}{dT_e} =$

$$r_1 + r_2 T_e + r_3$$

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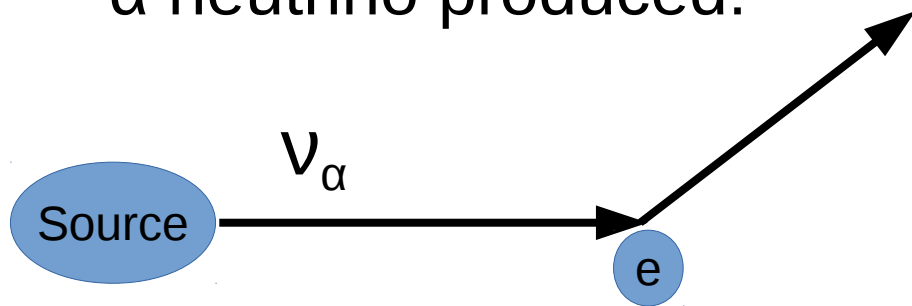
Current limits

- GEMMA, $\bar{\nu}_e$ from reactor: $|\mu_{\nu_e}| < 3 \times 10^{-11} \mu_B$.
- LSND, $\bar{\nu}_\mu, \nu_\mu$ from accelerator: $|\mu_{\nu_\mu}| < 7 \times 10^{-10} \mu_B$.
- DONUT, $\bar{\nu}_\tau, \nu_\tau$ from accelerator: $|\mu_{\nu_\tau}| < 4 \times 10^{-7} \mu_B$.
- GEMMA-II will improve by factor 3, SHiP could test nu-tau.
- Far from neutrino-induced $10^{-19} \mu_B$.
- (Borexino, solar ν , $3 \times 10^{-11} \mu_B$, see [talk by Oleg Smirnov](#).)

But what are we measuring here?

Effective magnetic moment $\mu_{\nu\alpha}$

- α neutrino produced:

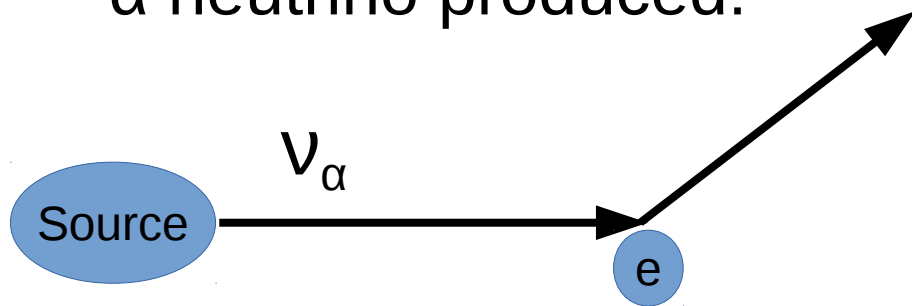


$\nu_{e,\mu,\tau}$ unobserved

$$\sigma \propto |\mu_{\nu\alpha}|^2 \equiv |\mu_{\alpha e}|^2 + |\mu_{\alpha\mu}|^2 + |\mu_{\alpha\tau}|^2$$

Effective magnetic moment $\mu_{\nu\alpha}$

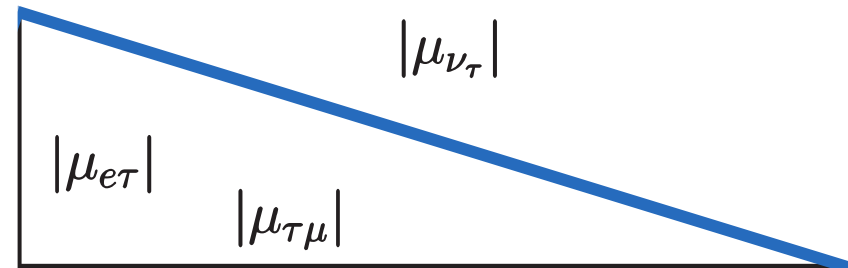
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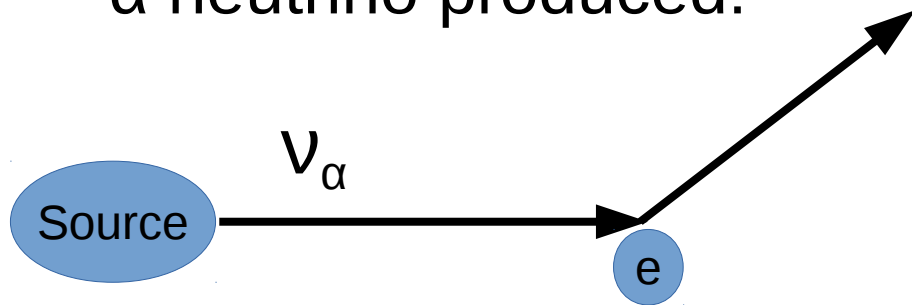
- For Majorana: $|\mu_{\alpha\alpha}| = 0$.



$$\mu_{\text{Maj.}} = \begin{pmatrix} 0 & \mu_{e\mu} & \mu_{e\tau} \\ -\mu_{e\mu} & 0 & \mu_{\mu\tau} \\ -\mu_{e\tau} & -\mu_{\mu\tau} & 0 \end{pmatrix}.$$

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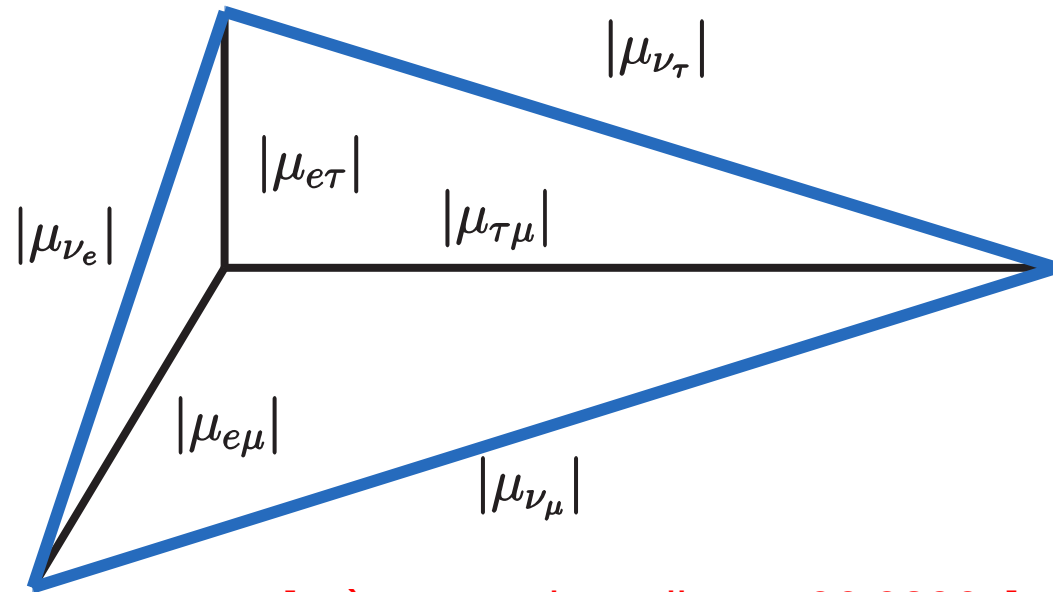
- Majorana triangle:

$$|\mu_{\nu_I}| \geq ||\mu_{\nu_J}| - |\mu_{\nu_K}||,$$

$$|\mu_{\nu_I}| \leq |\mu_{\nu_J}| + |\mu_{\nu_K}|.$$

- Special triangle:

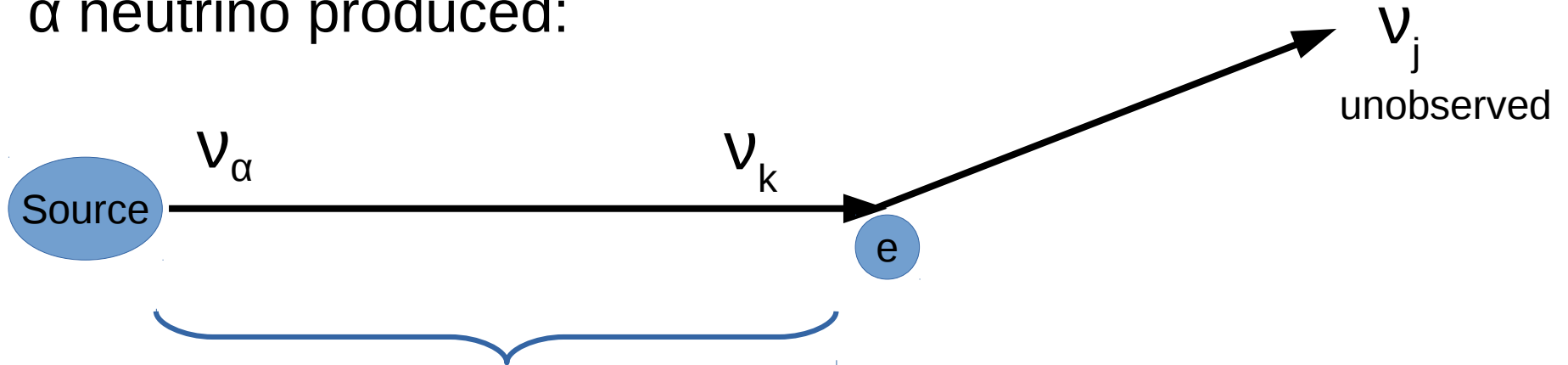
$$|\mu_{\nu_\tau}|^2 \leq |\mu_{\nu_\mu}|^2 + |\mu_{\nu_e}|^2 \text{ plus cyclic permutations.}$$



[Frère, Heeck, Mollet, 1506.02964]

Including neutrino mixing

- α neutrino produced:



- Oscillation length L between source and detector.
- Oscillation, then scattering into all mass eigenstates j :

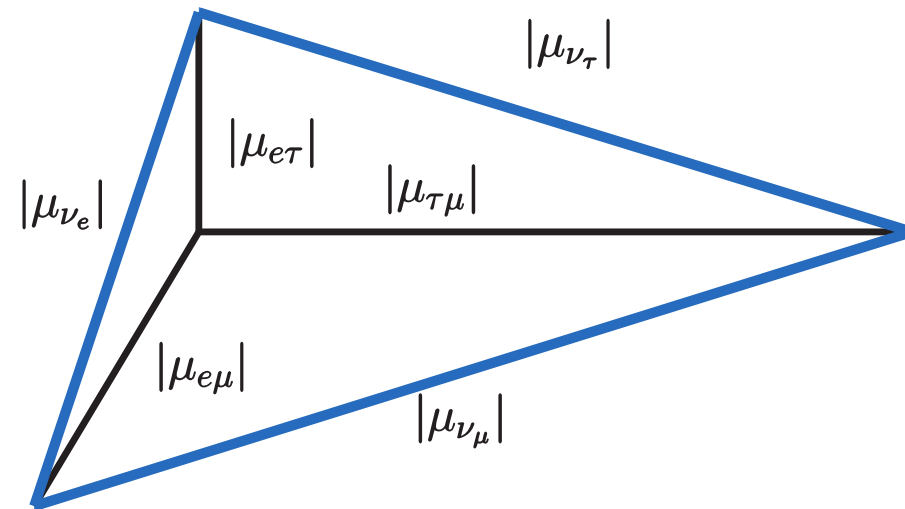
$$|\mu_{\nu_\alpha}|^2 = \sum_j \left| \sum_k U_{\alpha k}^* e^{-i\Delta m_{kj}^2 L / 2E_\nu} \mu_{jk} \right|^2 .$$

- In above experiments: $L/E \rightarrow 0$.

[Grimus, Stockinger, '98]

Triangle inequalities

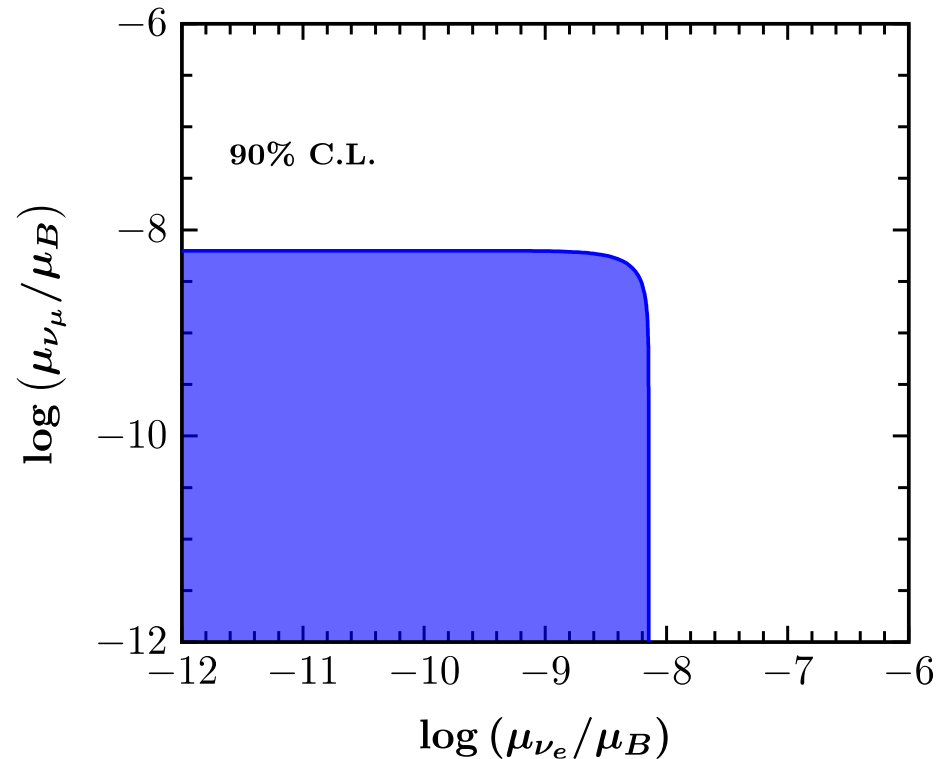
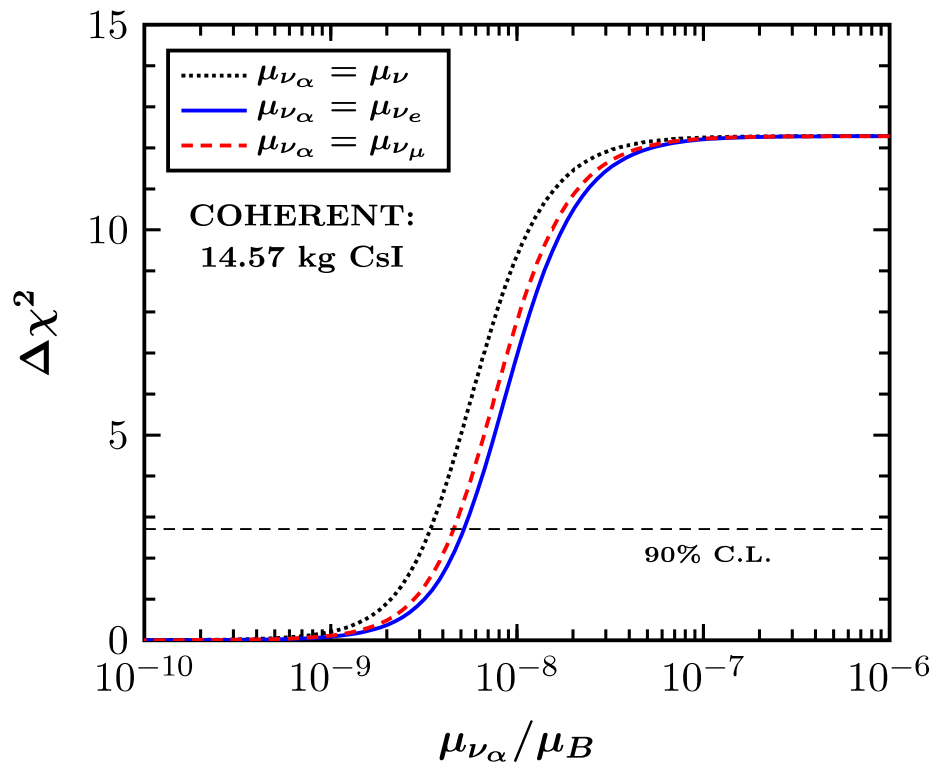
- Short distance $L/E \rightarrow 0$: $|\mu_{\nu_\alpha}|^2 = \sum_j |\sum_k U_{\alpha k}^* \mu_{jk}|^2$.
- For Majorana ν : μ antisymmetric!
- For unitary U , same inequalities as before,
 $|\mu_{\nu_\tau}|^2 \leq |\mu_{\nu_\mu}|^2 + |\mu_{\nu_e}|^2$ plus cyclic permutations.
- If violated \Rightarrow Not 3 Majorana ν !
- E.g. by SHiP measuring $\mu_{\nu\tau}$.
- Implies Dirac or light sterile ν .



[Frère, Heeck, Mollet, 1506.02964]

COHERENT

- Coherent elastic ν -nucleus scattering: $\frac{d\sigma}{dT_N} \propto \mu^2 \frac{Z^2}{T_N}$
[Dodd, Papageorgiu, Ranfone, '91; Kosmas+, '15]
- Pion at rest: mixture of ν_e and ν_μ .
- Improvement possible!



[Kosmas, Papoulias, 1711.09773]

How to get large μ_ν ?

- Main problem: $\mu_\nu \propto \delta m_\nu / \Lambda^2$.

[Voloshin, '88; Davidson+, '05; Bell+, '06]

- Light new physics?

$$\rightarrow \mu_\nu / \delta m_\nu = \epsilon e / 4M^2.$$

- Majorana ν can have

$$\mu_\nu \sim 10^{-12} \mu_B:$$

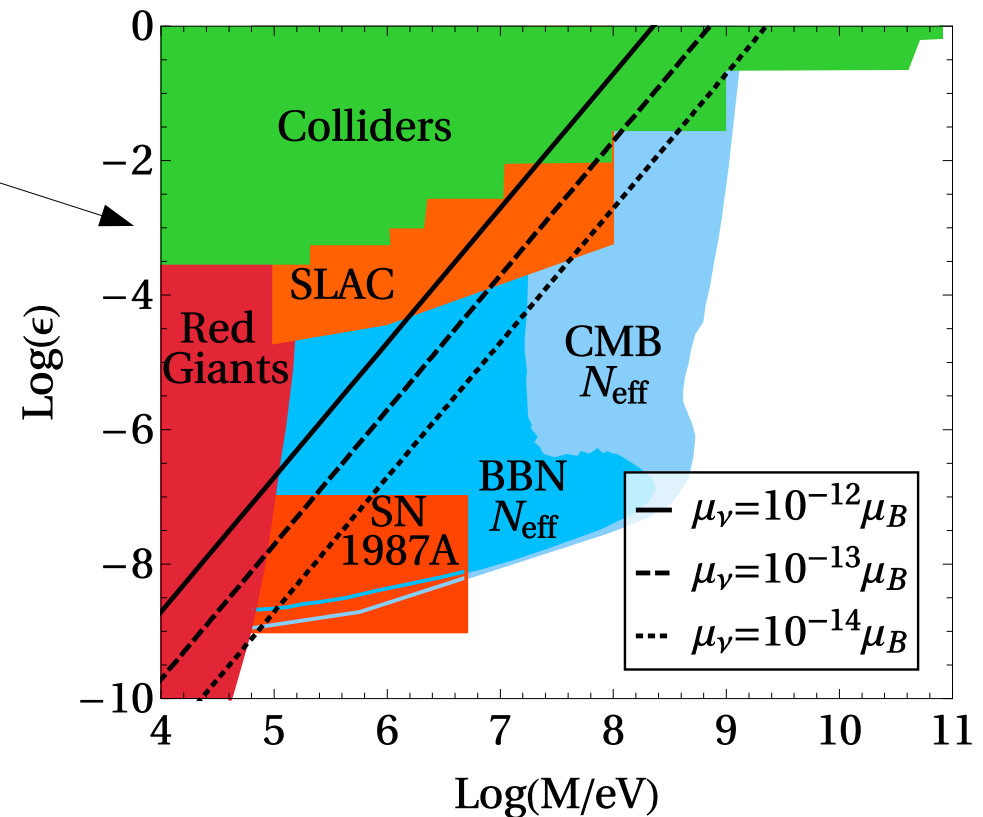
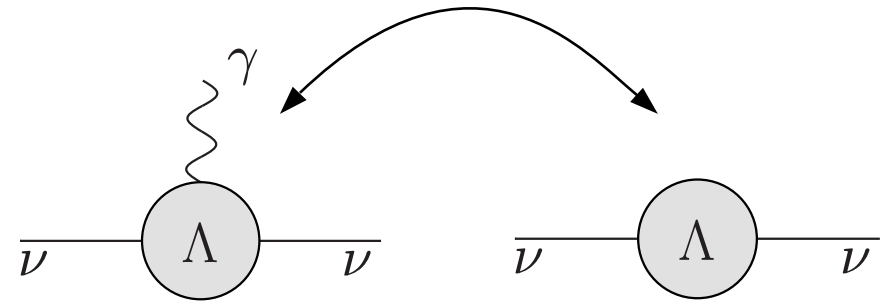
- Horizontal $SU(2)_H$.

[Babu, Mohapatra, '89]

- Barr-Freire-Zee model.

[Barr, Freire, Zee, '90]

- Dirac ν : need finetuning!



[Lindner, Radovčić, Welter, 1706.02555]

Barr-Freire-Zee model

- Spin-suppression:

$$\delta m_{ab} \propto \frac{m_b^2 - m_a^2}{M_W^2} \mu_{ab} M_{\text{scalar}}^2.$$

- Zee model:
3 doublets φ_a ,
1 singlet h^+ :

$$f_{ab} h^+ L_a L_b + M_{ab} h^+ \varphi_a \varphi_b + \lambda_{ab} \varphi_1^\dagger \varphi_a \varphi_1^\dagger \varphi_b.$$

- Majorana ν mass at loop level.
- Scalars @ TeV \Rightarrow still easily $\mu_\nu \sim 10^{-12} \mu_B$.

[Lindner, Radovčić, Welter, 1706.02555]

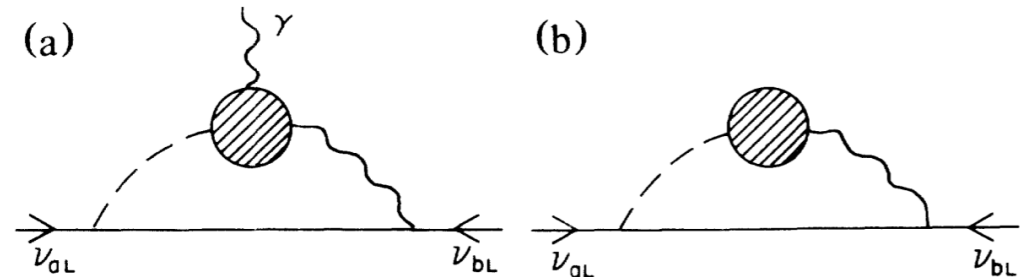


FIG. 1. (a) A graph that gives a contribution to the neutrino magnetic dipole moment. (b) The same graph with the photon removed contributes to the neutrino mass. However, this effect is suppressed by two powers of lepton mass because the virtual vector must be longitudinally polarized.

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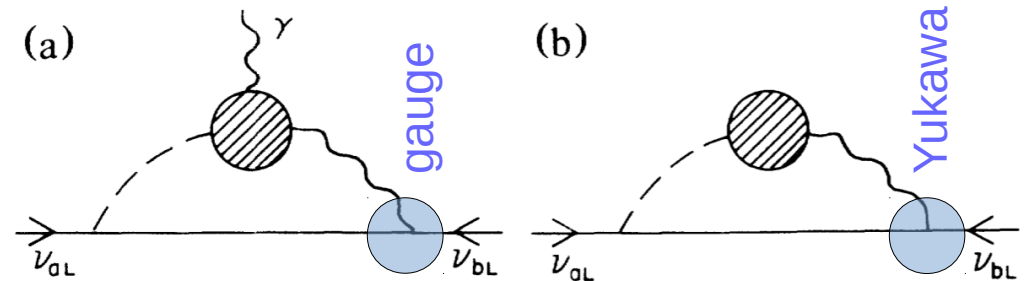


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[Barr, Freire, Zee, '90]

Summary

- m_ν induced magnetic moment: $\mu_\nu < 10^{-19} \mu_B$.
- Astrophysics reaches $10^{-12} \mu_B$, lab $10^{-11} \mu_B$.
- Improvement with GEMMA, COHERENT, SHiP,...
- Difficult to distinguish Majorana vs. Dirac.
- Model-building required for testable μ_ν .

Neutrinos always good for a surprise!